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# Positive solutions with a complex behavior for superlinear indefinite ODEs on the real line 

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#### Abstract

We show the existence of infinitely many positive solutions, defined on the real line, for the nonlinear scalar ODE $$
\ddot{u}+\left(a^{+}(t)-\mu a^{-}(t)\right) u^{3}=0
$$ where $a$ is a periodic, sign-changing function, and the parameter $\mu>0$ is large. Such solutions are characterized by the fact of being either small or large in each interval of positivity of $a$. In this way, we find periodic solutions, having minimal period arbitrarily large, and bounded non-periodic solutions, exhibiting a complex behavior. The proof is variational, exploiting suitable natural constraints of Nehari type.


AMS-Subject Classification. 34B18, 34C25, 34C28, 37J45
Keywords. Periodic and subharmonic solutions, natural constraints, Nehari method, singularly perturbed problems

## 1 Introduction

In this paper, we deal with the existence of positive bounded solutions, with a complex behavior, of the nonlinear scalar ODE

$$
\begin{equation*}
\ddot{u}+q(t) u^{3}=0, \quad t \in \mathbb{R} \tag{1}
\end{equation*}
$$

where $q(t)$ is a bounded and $T$-periodic function (for some $T>0$ ) which changes its sign. According to a terminology which is now quite standard in this setting (see [20]), equation (1) is thus superlinear indefinite.

When understanding $t$ as a space variable, equation (1) can be seen as a toy model of the elliptic PDE

$$
\begin{equation*}
\Delta u+\lambda u+q(x) u^{p}=0, \quad x \in \Omega \subset \mathbb{R}^{d} \tag{2}
\end{equation*}
$$

with $\lambda \in \mathbb{R}$ and $p>1$, which in turns arises when searching for steady states of the corresponding evolutionary parabolic problem (see [1] for a recent survey on the topic). Such kind of equations has a typical interpretation in the context of population dynamics, with the unknown $u$ playing the role of density of a species inhabiting the spatially heterogeneous domain

[^0]$\Omega$. Accordingly, the (indefinite) sign of the coefficient $q$ expresses saturation or autocatalytic behavior of the species $u$, when $q \leq 0$ or $q \geq 0$ respectively. Classical existence results, obtained both with topological and variational methods, for positive solutions of boundary value problems associated with (2) can be found among others in $[2,3,4,5,15,21]$.

Wishing to investigate the complexity of the solution set for (2), a typical strategy requires to play with the nodal behavior of the weight function $q$ and/or - in the PDE case - with the shape of $\Omega$, so as to regard (1)-(2) as singular perturbation problems (see [10, 11, 19]). For instance, in $[14,16]$ Gaudenzi, Habets and Zanolin dealt with the ODE (1) assuming that

$$
q(t)=a_{\mu}(t):=a^{+}(t)-\mu a^{-}(t),
$$

with $\mu$ a real parameter and $a^{+}, a^{-}$the positive and the negative part of a sign-changing function $a$, and they proved the existence of multiple positive solutions for

$$
\begin{equation*}
\ddot{u}+a_{\mu}(t) u^{3}=0 \tag{3}
\end{equation*}
$$

when $\mu \gg 1$. More precisely, they showed that the two-point boundary value problem $u(0)=u(T)=0$ associated with (3) has at least $2^{n}-1$ positive solutions (for $\mu$ large) whenever the weight function $a$ has $n$ disjoint intervals of positivity in $[0, T]$ (separated by intervals of negativity). The number $2^{n}-1$ comes from the possibility of prescribing, for a positive solution of (3), the behavior on each interval of positivity of $a$ among two possible ones: either the solution is "small" or is "large" (notice that the solution small on all the intervals is excluded, since it corresponds to the trivial one). Such a result, which was originally proved with a shooting technique, has later been generalized, using variational tools, by Bonheure, Gomes and Habets [6] and by Girão and Gomes [17, 18] to the Dirichlet problem for the corresponding elliptic PDE. Very recently, a topological approach for the ODE case has also been proposed by Feltrin and Zanolin [13].

It is the aim of the present paper to show that such kind of results has a natural analogue when the equation (3) is considered on an infinite interval (with $a$ a $T$-periodic function). That is, we can still produce, for $\mu$ sufficiently large, positive solutions of (3) defined on the real line and being either small or large on the intervals of positivity of the weight function $a$ according to a prescribed rule (which now involves the behavior of solutions on infinitely many intervals). More precisely, we prove the existence of positive $m T$-periodic solutions to (3) for any integer $m \geq 1$ (namely, positive subharmonic solutions) and, eventually, of bounded non-periodic solutions with a complex behavior, according to a typical scheme of "chaotic dynamics". Here is a simplified statement of our main existence result (for more details see Section 2).

Theorem 1.1. Let a be $L^{\infty}(\mathbb{R})$, T-periodic, and such that, for every $i \in \mathbb{Z}$,

$$
a(t) \geq 0, \quad \text { a.e. } \text { on } I_{i}^{+}, \quad a(t) \leq 0, \quad \text { a.e. on } I_{i}^{-}
$$

(but not identically zero) where the consecutive closed intervals $I_{i}^{+}, I_{i}^{-}$are such that $[0, T]$ is the union of a finite number of them.

For every integer $k \geq 1$, there exists $\mu^{*}>0$ such that, for every $\mu>\mu^{*}$ and for every double-sequence $\mathcal{L} \in\{0,1\}^{\mathbb{Z}}$ containing strings of zeroes of length at most $k$, the equation (3) has a positive solution $u \in W^{2, \infty}(\mathbb{R})$ such that, for every $i \in \mathbb{Z}$,

$$
\left.u\right|_{I_{i}^{+}} \quad \text { is "small" if } \mathcal{L}_{i}=0 \quad \text { and }\left.\quad u\right|_{I_{i}^{+}} \text {is "large" if } \mathcal{L}_{i}=1
$$

(and $\left.u\right|_{I_{i}^{-}}$is "small" for every $i$ ).
Moreover, such a solution can be chosen to be periodic whenever the sequence $\mathcal{L}$ is periodic.

It is worth mentioning that some results dealing with complex dynamics for ODEs with indefinite weight have already appeared (see, among others, [8, 9, 26, 28]). However, [9, 26, 28] deal with oscillatory solutions of the superlinear indefinite equation (1). On the other hand, a chaotic dynamics entirely made by positive solutions for an equation like $\ddot{u}+q(t) g(u)=0$ is produced in [8], but in the case of a nonlinearity $g(u)$ of super-sublinear type; the topological technique (the so-called "Stretching Along the Paths" method) used therein, moreover, is not applicable in the present context.

In our result, periodic solutions play a crucial role and, indeed, bounded non-periodic solutions are constructed as limit of periodic ones when the period goes to infinity. In doing this, a careful analysis has to be performed in order to ensure that the constant $\mu^{*}$ for which $m T$-periodic solutions are available can be chosen independently on the integer $m$, and thus making possible the passage to the limit $m \rightarrow+\infty$. We stress that this procedure (obtaining chaotic solutions as limit of periodic solutions) does not seem to be easily reproducible using the approaches proposed in $[6,13,14,16,17,18]$.

Our proof of the existence of periodic solutions is variational: we exploit the fact that solutions of (3) are critical points of the related action functional $J$, and we construct a suitable natural constraint for such a functional, that is, a constraint for which constrained critical points of $J$ are free ones. The most famous natural constraint is the Nehari manifold, which can be successfully used in order to find positive solutions of the two-point boundary value problem for (3) when restricted to intervals where $a$ is non negative. Indeed, as shown in [22, 23], letting

$$
\mathcal{N}_{i}=\left\{u \in H_{0}^{1}\left(I_{i}^{+}\right): u \not \equiv 0, \int_{I_{i}^{+}} \dot{u}^{2}=\int_{I_{i}^{+}} a^{+} u^{4}\right\} \quad \text { and } \quad c_{i}=\inf _{u \in \mathcal{N}} \frac{1}{4} \int_{I_{i}^{+}} \dot{u}^{2}
$$

we have that the set

$$
\begin{equation*}
\mathcal{K}_{i}=\left\{u \in \mathcal{N}_{i}: \frac{1}{4} \int_{I_{i}^{+}} \dot{u}^{2}=c_{i}\right\} \subset W^{2, \infty}\left(I_{i}^{+}\right) \tag{4}
\end{equation*}
$$

consists of one-sign solutions of (3) with homogeneous Dirichlet boundary conditions on $I_{i}^{+}$. Under this perspective, one may read Theorem 1.1 as a singular perturbation result, where the singular limit of the solutions we find, as $\mu \rightarrow+\infty$, is the set

$$
\begin{equation*}
\mathcal{K}_{\mathcal{L}}=\left\{u \in W^{1, \infty}(\mathbb{R}):\left.u\right|_{I_{i}^{+}} \in \mathcal{K}_{i},\left.u\right|_{I_{i}^{+}}>0 \text { if } \mathcal{L}_{i}=1 \text { and } u \equiv 0 \text { elsewhere }\right\} . \tag{5}
\end{equation*}
$$

In fact, we consider a periodic truncation of such a set, and we deform it to a suitable Neharitype constraint, showing that minima of the action functional on such a set, when $\mu$ is large, correspond to the desired periodic orbits. The same idea was already exploited in [6, 17, 18], even though such papers concern the PDE setting, with Dirichlet boundary conditions on bounded domains, and it is not clear how to modify the arguments there, in order to treat periodic conditions, and to obtain uniform estimates so that one can pass to the limit to unbounded domains. To overcome this difficulty, we rely on an abstract result contained in [24]: this avoids the necessity of constructing a projection operator to the constraint, which is usually one of the most delicate parts when dealing with Nehari-type arguments. The pay-off of such a method is that it provides with sharp localization of the solutions, and it allows to prove optimal bounds, uniform as $\mu \rightarrow+\infty$.

Theorem 1.2. In the assumptions of Theorem 1.1, let $\mathcal{L}$ be fixed, and for every $\mu>\mu^{*}$ let $u_{\mu}$ denote the corresponding solution. Then there exists a constant $C=C(a, \mathcal{L})$, not depending on $\mu$, such that

$$
\left\|u_{\mu}\right\|_{C^{0,1}(\mathbb{R})}<C
$$

Furthermore, up to subsequences,

$$
u_{\mu} \rightarrow \bar{u} \in \mathcal{K}_{\mathcal{L}}, \quad \text { in } C^{0, \alpha}(\mathbb{R}), \text { for every } \alpha<1
$$

and the convergence is true also in $H_{\mathrm{loc}}^{1}(\mathbb{R})$, and in $W_{\mathrm{loc}}^{2, \infty}$ away from the points where the function a changes sign.

As we mentioned, the bounds above are optimal: since $u_{\mu}$ is $C^{1}$, were the convergence $C^{0,1}$, also $\bar{u}$ would be $C^{1}$; this is impossible, since the elements of $\mathcal{K}_{\mathcal{L}}$ cannot be $C^{1}$ (if $\mathcal{L}_{i} \neq 0$ for at least one $i$ ).

Variational methods were already successfully exploited to construct entire complex solutions of nonlinear oscillators in [28, 30, 25, 27]. Though also the methods employed in these papers are related to the results by Nehari, in particular to a broken geodesics argument, that situation is rather different: the solutions found there are oscillatory, and the uniform energy estimates to pass from bounded intervals to the real line are obtained through a control of the distance between consecutive zeroes. Rather, we borrow some ideas from [29], where radial positive multi-bump solutions to a singularly perturbed elliptic system are investigated.

To conclude, we remark that slight variants of our technique can be invoked to prove some related results (see Remark 7.1 at the end of the paper). First of all, one can consider functions $a$ changing sign with some uniform properties, rather than periodic ones; further, also changing sign solutions can be constructed, by choosing sequences $\mathcal{L} \in\{-1,0,1\}^{\mathbb{Z}}$ and prescribing $\left.u\right|_{I_{i}^{+}}$to be large and positive (resp. negative) whenever $\mathcal{L}_{i}=1$ (resp. -1 ). Moreover, we can prove the existence of $2^{n}-1$ positive solutions (or $3^{n}-1$ nontrivial ones, possibly changing sign) to the $T$-periodic boundary value problem (that is, $u(0)=u(T)$ and $\dot{u}(0)=\dot{u}(T))$ associated with (3), whenever the function $a$ has $n$ intervals of positivity in a period. This gives a $T$-periodic counterpart of the result first proved by Gaudenzi, Habets and Zanolin for the Dirichlet problem. The Neumann boundary value problem $\dot{u}(0)=\dot{u}(T)=0$ could be also successfully considered, using very similar arguments, and thus extending [7]. We stress, however, that all these results, dealing with a boundary value problem on a finite interval, can be obtained with much easier arguments (on the lines of the main application in [24]) with respect to the ones described in this paper, whose crucial theme insists on finding estimates for the threshold $\mu^{*}$, which are independent of the size of the considered interval.

## 2 Main result and strategy of the proof

In this paper we deal with the existence of positive solutions of the superlinear indefinite equation

$$
\begin{equation*}
\ddot{u}+a_{\mu}(t) u^{3}=0, \quad t \in \mathbb{R}, \tag{6}
\end{equation*}
$$

where $\mu>0$ is a large parameter and

$$
a_{\mu}(t+T)=a_{\mu}(t):=a^{+}(t)-\mu a^{-}(t), \quad \text { for every } t
$$

with $a^{+}(t)=\max (0, a(t))$ and $a^{-}(t)=\max (0,-a(t))$ denoting the positive and negative part of a sign-changing, $T$-periodic function $a \in L^{\infty}(\mathbb{R})$. For the sake of simplicity, we assume that $a$ changes sign just once in $[0, T]$, that is:
(A) there exists $\tau \in] 0, T[$ such that

$$
a(t) \geq 0, \not \equiv 0 \text { on }[0, \tau], \quad a(t) \leq 0, \not \equiv 0 \text { on }[\tau, T]
$$

(even though we can treat much more general situations, see Remark 7.1 at the end of the paper). Up to a time-translation and a suitable choice of $\tau$ we can suppose that

$$
\begin{equation*}
\int_{\tau}^{\tau+\delta} a^{-}(t) d t>0 \quad \text { and } \quad \int_{T-\delta}^{T} a^{-}(t) d t>0 \tag{7}
\end{equation*}
$$

for every small $\delta>0$.
From now on, we also use the notation

$$
\sigma_{i}=i T, \quad \tau_{i}=\tau+i T, \quad I_{i}^{+}=\left[\sigma_{i}, \tau_{i}\right], \quad I_{i}^{-}=\left[\tau_{i}, \sigma_{i+1}\right], \quad \text { for every } i \in \mathbb{Z}
$$

Our main result reads as follows. In the statement below, $\{0,1\}^{\mathbb{Z}}$ denotes the space of double-sequences of 0 and 1 . Moreover, for $\mathcal{L}=\left\{\mathcal{L}_{i}\right\}_{i \in \mathbb{Z}} \in\{0,1\}^{\mathbb{Z}}$, we set

$$
\mathbf{0}_{\mathcal{L}}=\sup \left\{n \in \mathbb{N}: \exists i \in \mathbb{Z} \text { s.t. } \mathcal{L}_{j}=0, \forall j=i, \ldots, i+n-1\right\}
$$

namely, the maximal length of strings in $\mathcal{L}$ entirely composed by 0 .
Theorem 2.1. For every integer $k \geq 1$ there exists $\mu^{*}=\mu^{*}(k)>0$ such that, for every $\mu>\mu^{*}$ and for every $\mathcal{L} \in\{0,1\}^{\mathbb{Z}}$ with

$$
\begin{equation*}
\boldsymbol{O}_{\mathcal{L}} \leq k \tag{8}
\end{equation*}
$$

equation (6) has a positive solution $u \in W^{2, \infty}(\mathbb{R})$ such that, for every $i \in \mathbb{Z}$,

$$
\begin{equation*}
\int_{I_{i}^{+}} \dot{u}^{2}<r^{2} \text { if } \quad \mathcal{L}_{i}=0 \quad \text { and } \quad \int_{I_{i}^{+}} \dot{u}^{2}>r^{2} \text { if } \mathcal{L}_{i}=1, \tag{9}
\end{equation*}
$$

and

$$
\begin{equation*}
\|\dot{u}\|_{L^{\infty}(\mathbb{R})} \leq \rho, \tag{10}
\end{equation*}
$$

where $r$ and $\rho$ are positive explicit constants, only depending on the weight a.
More precisely: for any $\varepsilon>0$ there exists $\mu^{*}=\mu^{*}(k, \varepsilon) \geq \mu^{*}(k)$, with $\mu^{*}(k, \varepsilon) \rightarrow+\infty$ for $\varepsilon \rightarrow 0^{+}$, such that the solution $u$ fulfills, for any $i \in \mathbb{Z}$,
(P1) $\|u\|_{L^{\infty}\left(I_{i}^{-}\right)}+\int_{I_{i}^{-}} \dot{u}^{2} \leq \varepsilon$,
(P2) $\|u\|_{W^{2, \infty}\left(I_{i}^{+}\right)} \leq \varepsilon$ if $\mathcal{L}_{i}=0$,
(P3) $\operatorname{dist}_{W^{2, \infty}}\left(\left.u\right|_{I_{i}^{+}}, \mathcal{K}_{i}\right) \leq \varepsilon$ if $\mathcal{L}_{i}=1$, where $\mathcal{K}_{i}$ is defined in (4).
Finally, the solution can be chosen to be mT-periodic whenever the sequence $\mathcal{L}$ is $m$ periodic for some $m \in \mathbb{N}$.

The proof of Theorem 2.1 is based on an approximation procedure which we now describe. For each integer $N \geq 0$, consider the interval $I_{N}=\left[\sigma_{-N}, \sigma_{N+1}\right]$, that is,

$$
I_{N}=I_{-N}^{+} \cup I_{-N}^{-} \cup \ldots \cup I_{N}^{+} \cup I_{N}^{-}
$$

Then, the following result holds true.
Theorem 2.2. For any integer $k \geq 1$, there exists $\mu^{*}>0$ such that for every $\mu>\mu^{*}$, $\mathcal{L} \in\{0,1\}^{\mathbb{Z}}$ satisfying (8) and $N>k$, equation

$$
\begin{equation*}
\ddot{u}+a_{\mu}(t) u^{3}=0, \quad t \in I_{N} \tag{11}
\end{equation*}
$$

has a positive solution $u \in W_{\text {per }}^{2, \infty}\left(I_{N}\right)$ such that properties (9), (10), (P1), (P2) and (P3) hold, for $i=-N, \ldots, N$.

Notice that, whenever the sequence $\mathcal{L}$ in Theorem 2.1 is $m$-periodic, with $m$ an odd integer number, Theorem 2.2 gives the existence of a positive $m T$-periodic solution to (6). The case in which $m$ is an even integer can be handled in a completely analogous way by considering, for $N \geq 1$, the interval $\widetilde{I}_{N}=\left[\sigma_{-N}, \sigma_{N-1}\right]$.

On the other hand, for non-periodic sequences $\mathcal{L}$, the corresponding positive solution $u=u_{\mathcal{L}} \in W^{2, \infty}(\mathbb{R})$ of Theorem 2.1 can be constructed as limit, for $N \rightarrow+\infty$, of the solutions $u_{\mathcal{L}, N} \in W_{\text {per }}^{2, \infty}\left(I_{N}\right)$ found in Theorem 2.2. More details for this (quite standard) argument will be given at the end of the paper, in Section 7.

From now on, we will concentrate on Theorem 2.2, whose proof will take a great part of the paper. It relies on a variational argument, consisting in the minimization of the action functional

$$
J_{\mu, I_{N}}(u)=\frac{1}{2} \int_{I_{N}} \dot{u}^{2}-\frac{1}{4} \int_{I_{N}} a_{\mu} u^{4}
$$

on suitable Nehari-type subset of $H_{\text {per }}^{1}\left(I_{N}\right)$ (in the following, not to overload the notation, we will often drop the subscript $I_{N}$ when no confusion is possible). To describe our procedure, we need some preliminary notation.

First of all, we define a cut-off function $\eta \in C_{c}^{\infty}(\mathbb{R})$ such that $0 \leq \eta(t) \leq 1$, for every $t \in \mathbb{R}$ and

$$
\eta \equiv 1 \text { on }[0, \tau] \quad \text { and } \quad \eta \equiv 0 \text { on } \mathbb{R} \backslash\left[\frac{\tau-T}{4}, \tau+\frac{T-\tau}{4}\right]
$$

Moreover, we set

$$
\eta_{i}(t)=\eta\left(t-\sigma_{i}\right), \quad \text { for every } i \in \mathbb{Z}
$$

so as to obtain a family of cut-off functions $\left\{\eta_{i}\right\}_{i \in \mathbb{Z}}$ such that $\eta_{i} \eta_{j} \equiv 0$ whenever $i \neq j$.
Next, we turn to introduce the Nehari-type constraint, which depends on $N$ and $\mathcal{L}$. Setting

$$
L=\left\{i \in\{-N, \ldots, N\}: \mathcal{L}_{i}=1\right\}
$$

we define the subspaces

$$
V^{+}=\left\{u \in H_{\mathrm{per}}^{1}\left(I_{N}\right): u \equiv 0 \text { on } \bigcup_{i \in L} I_{i}^{+}\right\}
$$

and, for any $u \in H_{\text {per }}^{1}\left(I_{N}\right)$,

$$
V_{u}^{-}=\operatorname{span}_{i \in L}\left\{\eta_{i} u\right\}
$$

The Nehari-type set is then defined as

$$
\begin{equation*}
\mathcal{P}_{\mu}=\left\{u \in H_{\mathrm{per}}^{1}\left(I_{N}\right): \operatorname{proj}_{\left(V^{+} \oplus V_{u}^{-}\right)} \nabla J_{\mu}(u)=0\right\} . \tag{12}
\end{equation*}
$$

The next result collects some properties enjoyed by the functions in $\mathcal{P}_{\mu}$.
Lemma 2.3. Let $u \in \mathcal{P}_{\mu}$. Then (all integrals are understood on $I_{N}$ ):
(i) $\int \dot{u} \dot{v}=\int a_{\mu} u^{3} v$ for every $v \in V^{+}$. In particular, $u \in W^{2, \infty}\left(I_{N} \backslash \cup_{i \in L} I_{i}^{+}\right)$and

$$
\ddot{u}+a_{\mu}(t) u^{3}=0, \quad \text { on } I_{N} \backslash \cup_{i \in L} I_{i}^{+} .
$$

(ii) $\int \dot{u}\left(\dot{\eta_{i}} u\right)=\int a_{\mu} \eta_{i} u^{4}$ for every $i \in L$; equivalently,

$$
\int_{I_{i}^{+}}\left(\dot{u}^{2}-a^{+} u^{4}\right)=u\left(\tau_{i}\right) \dot{u}\left(\tau_{i}^{+}\right)-u\left(\sigma_{i}\right) \dot{u}\left(\sigma_{i}^{-}\right) .
$$

Here, $\dot{u}\left(\sigma_{i}^{-}\right)$and $\dot{u}\left(\tau_{i}^{+}\right)$are respectively the left derivative of $u$ in $\sigma_{i}$ and the right derivative of $u$ in $\tau_{i}$, whose existence is guaranteed by the previous point (i).
(iii) $\int \dot{u}^{2}=\int a_{\mu} u^{4}$; hence

$$
J_{\mu}(u)=\frac{1}{4} \int_{I_{N}} \dot{u}^{2}=\frac{1}{4} \int_{I_{N}} a_{\mu} u^{4} .
$$

(iv) $\int \dot{u}\left(\eta_{i}^{2} u\right)=\int a_{\mu} \eta_{i}^{2} u^{4}$ for every $i=-N, \ldots, N$; as a consequence

$$
\int_{I_{N}}\left(\eta_{i} u\right)^{2}=\int_{I_{N}} a_{\mu} \eta_{i}^{2} u^{4}+\int_{I_{N}}{\dot{\eta_{i}}}^{2} u^{2}
$$

Proof. We prove separately each point.
(i) This corresponds to $\operatorname{proj}_{V^{+}} \nabla J_{\mu}(u)=0$. In particular, this implies that $u$ solves $\ddot{u}+$ $a_{\mu}(t) u^{3}=0$ in the sense of distributions on $I_{N} \backslash \cup_{i \in L} I_{i}^{+}$, and the second claim follows by elliptic regularity.
(ii) The first equality corresponds to $\operatorname{proj}_{V_{u}^{-}} \nabla J_{\mu}(u)=0$. Denoting $\operatorname{supp}\left(\eta_{i}\right)=\left[\sigma_{i}^{\prime}, \tau_{i}^{\prime}\right]$, such an equality writes as

$$
\begin{aligned}
\int_{I_{i}^{+}}\left(\dot{u}^{2}-a^{+} u^{4}\right) & =-\int_{\left[\sigma_{i}^{\prime}, \sigma_{i}\right] \cup\left[\tau_{i}, \tau_{i}^{\prime}\right]} \dot{u}\left(\dot{\eta_{i}} u\right)-a_{\mu} \eta_{i} u^{4} \\
& =-\int_{\left[\sigma_{i}^{\prime}, \sigma_{i}\right] \cup\left[\tau_{i}, \tau_{i}^{\prime}\right]}\left(-\ddot{u}-a_{\mu} u^{3}\right)\left(\eta_{i} u\right)+u\left(\tau_{i}\right) \dot{u}\left(\tau_{i}^{+}\right)-u\left(\sigma_{i}\right) \dot{u}\left(\sigma_{i}^{-}\right),
\end{aligned}
$$

and the first term vanishes by (i).
(iii) The identity is equivalent to $\left\langle\nabla J_{\mu}(u), u\right\rangle=0$, then we just need to show that $u \in$ $V^{+} \oplus V_{u}^{-}$. Let us write $u \in H_{\text {per }}^{1}\left(I_{N}\right)$ as

$$
u=\sum_{i \in L} \eta_{i} u+\left(1-\sum_{i \in L} \eta_{i}\right) u
$$

then the first term lies in $V_{u}^{-}$, the second one in $V^{+}$.
(iv) Similarly we need to prove that $\eta_{i}^{2} u \in V^{+} \oplus V_{u}^{-}$, for every $i=-N, \ldots, N$. When $i \notin L$, we immediately conclude that $\eta_{i}^{2} u \in V^{+}$. On the other hand, when $i \in L$ we write

$$
\eta_{i}^{2} u=\eta_{i}\left(\eta_{i}-1\right) u+\eta_{i} u ;
$$

the first term belongs to $V^{+}$, while the second one to $V_{u}^{-}$.

Notice that, in view of Lemma 2.3 (iii), the functional $J_{\mu}$ is bounded below on $\mathcal{P}_{\mu}$. Our minimization problem will be settled in an open subset $\mathcal{N}_{\mu} \subset \mathcal{P}_{\mu}$. To describe it, we need to introduce some notation. First of all, define for some suitable $\zeta \in(0,(T-\tau) / 2)$

$$
\begin{align*}
\mathcal{N} & =\left\{u \in H_{0}^{1}(0, \tau): u \not \equiv 0, \int_{0}^{\tau} \dot{u}^{2}=\int_{0}^{\tau} a^{+} u^{4}\right\} \\
\mathcal{Z}_{\zeta} & =\{u \in \mathcal{N}: \exists \bar{t} \in[\zeta, \tau-\zeta] \text { s.t. } u(\bar{t})=0\} \tag{13}
\end{align*}
$$

and set

$$
\begin{equation*}
c=\inf _{u \in \mathcal{N}} \frac{1}{4} \int_{0}^{\tau} \dot{u}^{2} \quad \text { and } \quad c_{\zeta}=\inf _{u \in \mathcal{Z}_{\zeta}} \frac{1}{4} \int_{0}^{\tau} \dot{u}^{2} \tag{14}
\end{equation*}
$$

We claim here that

$$
c<c_{\zeta} .
$$

Indeed, notice first that $c \leq c_{\zeta}$ since $\mathcal{Z}_{\zeta} \subset \mathcal{N}$. Then, observe that $c$ and $c_{\zeta}$ are both attained (as infimum value of the corresponding minimizing problems): this is well known for $c$, and the same proof works also for $c_{\zeta}$. Since functions attaining the value $c$ are (non-trivial) constant-sign $H_{0}^{1}$-solutions of $\ddot{u}+a^{+} u^{3}=0$ on $(0, \tau)$ [22, 23], which are not in $\mathcal{Z}_{\zeta}$, we conclude $c \neq c_{\zeta}$.

Then, for suitable constants $r, K, \rho>0$, consider the following conditions
(C1) $\int_{I_{i}^{+}} \dot{u}^{2}<r^{2}$ for $i \notin L$ and $r^{2}<\int_{I_{i}^{+}} \dot{u}^{2}<2\left(c+c_{\zeta}\right)$ for $i \in L$,
(C2) $u(t)>0$ for every $t \in \bigcup_{i \in L}\left[\sigma_{i}+\zeta, \tau_{i}-\zeta\right]$,
(C3) $|u(t)|<K$ for every $t \in \bigcup_{i=-N}^{N} I_{i}^{-}$,
(C4) for $i \in L$,

$$
\begin{array}{lllll}
\dot{u}\left(\sigma_{i}^{-}\right)<\rho & \text { if } u\left(\sigma_{i}\right) \geq 0 & \text { and } & \dot{u}\left(\sigma_{i}^{-}\right)>-\rho & \text { if } u\left(\sigma_{i}\right) \leq 0 \\
\dot{u}\left(\tau_{i}^{+}\right)>-\rho & \text { if } u\left(\tau_{i}\right) \geq 0 & \text { and } & \dot{u}\left(\tau_{i}^{+}\right)<\rho & \text { if } u\left(\tau_{i}\right) \leq 0
\end{array}
$$

and set

$$
\begin{equation*}
\mathcal{N}_{\mu}=\left\{u \in \mathcal{P}_{\mu}: u \text { satisfies }(\mathrm{C} 1),(\mathrm{C} 2),(\mathrm{C} 3),(\mathrm{C} 4)\right\} \tag{15}
\end{equation*}
$$

The precise value for $\zeta, r$ and $\rho$ will be given in (25), (28) and (32) respectively. As for $K$, the choice is more arbitrary, since (as it will be clear from the proof) any positive value for it works, up to enlarging $\mu^{*}$. However, a natural choice can be made by recalling (see, for instance, [14, Lemma 4.3]) that positive solutions to (11) are $L^{\infty}$ a-priori bounded, independently on both $\mu>0$ and $N$.

Remark 2.1. A few comments on the set $\mathcal{N}_{\mu}$ are now in order. First, the derivatives involved in (C4) are well defined since $u$ solves (11) on each $I_{i}^{-}$; furthermore such a condition is open in the $H^{1}$-topology. Second, the set $\mathcal{N}_{\mu}$ is non-empty since, recalling (5),

$$
\mathcal{K}_{\mathcal{L}} \subset \mathcal{N}_{\mu}
$$

provided $r^{2}<4 c$.
As a first step towards Theorem 2.2, we have the following result, whose proof will be given in Section 4, taking advantage of some technical lemmas developed in Section 3, as well as of the main result in [24].

Proposition 2.4. There exists $\mu^{*}>0$ (depending on the weight function $a$ and on the integer $k$, but not on $\mathcal{L}$ and $N$ ) such that, for any $\mu>\mu^{*}$, the set $\mathcal{N}_{\mu}$ is a $C^{1}$ embedded submanifold of $H_{\mathrm{per}}^{1}\left(I_{N}\right)$ such that any constrained Palais-Smale sequence is a free one. That is, if $\left(u_{n}\right) \subset \mathcal{N}_{\mu}$ is such that $J_{\mu}\left(u_{n}\right)$ is bounded and $\nabla_{\mathcal{N}_{\mu}} J_{\mu}\left(u_{n}\right) \rightarrow 0$, then $\nabla J_{\mu}\left(u_{n}\right) \rightarrow 0$ as well.

According to the above result, the argument leading to Theorem 2.2 now proceeds by exhibiting a bounded constrained Palais-Smale sequence $\left(u_{n}\right) \subset \mathcal{N}_{\mu}$ at level

$$
\begin{equation*}
\inf _{u \in \mathcal{N}_{\mu}} J_{\mu}(u) \tag{16}
\end{equation*}
$$

Indeed, Proposition 2.4 implies that this is a free bounded Palais-Smale sequence and, since the gradient of $J_{\mu}$ is a compact perturbation of an invertible operator, a (free) critical point for $J_{\mu}$ is obtained. This is a solution of (11) having - by construction - the desired complex behavior. Sections 5 and 6 will be devoted to this delicate argument, which requires a careful understanding of the behavior of $J_{\mu}$ near the boundary of $\mathcal{N}_{\mu}$.

To conclude the proof of Theorem 2.2, it will be then enough to prove that the solution found is positive, uniformly bounded in $C^{1}$, and it has the required properties (P1), (P2) and (P3) for $\mu \rightarrow+\infty$. This will be the goal of Section 7, containing also some more details for the limit $N \rightarrow+\infty$ leading to Theorem 2.1.

## 3 Some technical results

Throughout the paper, we will make often use of the following elementary inequality:

$$
\begin{equation*}
\|u\|_{L^{\infty}\left(s_{1}, s_{2}\right)} \leq \min _{\left[s_{1}, s_{2}\right]}|u|+\sqrt{s_{2}-s_{1}}\left(\int_{s_{1}}^{s_{2}} \dot{u}^{2}\right)^{1 / 2}, \quad \text { for every } u \in H^{1}\left(s_{1}, s_{2}\right) \tag{17}
\end{equation*}
$$

This is a simple consequence of the fundamental theorem of calculus, together with the Cauchy-Schwartz inequality. Notice also that, if $u$ vanishes somewhere on $\left[s_{1}, s_{2}\right]$, then from (17) we obtain the Sobolev-type inequality

$$
\begin{equation*}
\|u\|_{L^{\infty}\left(s_{1}, s_{2}\right)}^{2} \leq\left(s_{2}-s_{1}\right)\|\dot{u}\|_{L^{2}\left(s_{1}, s_{2}\right)}^{2} \tag{18}
\end{equation*}
$$

and the Poincaré-type inequality

$$
\begin{equation*}
\|u\|_{L^{2}\left(s_{1}, s_{2}\right)}^{2} \leq\left(s_{2}-s_{1}\right)^{2}\|\dot{u}\|_{L^{2}\left(s_{1}, s_{2}\right)}^{2} . \tag{19}
\end{equation*}
$$

### 3.1 Local estimates on the solutions

In this section, we collect some useful estimates for solutions of the differential equation (6) in an interval of positivity of the weight function, say

$$
\begin{equation*}
\ddot{u}+a^{+}(t) u^{3}=0, \quad t \in[0, \tau], \tag{20}
\end{equation*}
$$

and in an interval of negativity, say

$$
\begin{equation*}
\ddot{u}-\mu a^{-}(t) u^{3}=0, \quad t \in[\tau, T] . \tag{21}
\end{equation*}
$$

Obviously, in view of the $T$-periodicity of $a$, identical results will hold true for solutions on the intervals $I_{i}^{+}$and $I_{i}^{-}$for every $i \in \mathbb{Z}$. Let us also observe once for all that solutions $u$
of (20) are concave (resp., convex) when $u \geq 0$ (resp., $u \leq 0$ ), while solutions $u$ of (21) are convex (resp., concave) when $u \geq 0$ (resp., $u \leq 0$ ) and satisfy

$$
\begin{equation*}
|u(t)| \leq \max (|u(\tau)|,|u(T)|), \quad \text { for every } t \in[\tau, T] \tag{22}
\end{equation*}
$$

and

$$
\begin{equation*}
|\dot{u}(t)| \leq \max \left(\left|\dot{u}\left(\tau^{+}\right)\right|,\left|\dot{u}\left(T^{-}\right)\right|\right), \quad \text { for every } t \in[\tau, T] \tag{23}
\end{equation*}
$$

These facts will be used several times without further comments.
The first result is a $C^{1}$ a-priori estimate in $[0, \tau]$, given a bound for $\int_{0}^{\tau} \dot{u}^{2}$.
Lemma 3.1. For every $M>0$ there exists $M^{\prime}>0$ such that, for any u solving (20), it holds

$$
\int_{0}^{\tau} \dot{u}^{2} \leq M \quad \Longrightarrow \quad\|u\|_{L^{\infty}(0, \tau)}+\|\dot{u}\|_{L^{\infty}(0, \tau)} \leq M^{\prime}
$$

Proof. To see this, we first observe that it is enough to prove the boundedness of $\|u\|_{L^{\infty}(0, \tau)}$. Indeed, since $u$ solves (20), the boundedness of $\|\dot{u}\|_{L^{\infty}(0, \tau)}$ follows from the ones of both $u$ and $\ddot{u}$ in $L^{\infty}$, via the elementary inequality

$$
\begin{equation*}
\|\dot{u}\|_{L^{\infty}(0, \tau)} \leq \frac{2}{\tau}\|u\|_{L^{\infty}(0, \tau)}+\tau\|\ddot{u}\|_{L^{\infty}(0, \tau)} . \tag{24}
\end{equation*}
$$

In order to conclude, by virtue of (17), we just need to show that $\min _{[0, \tau]}|u|$ is bounded. When $u$ vanishes at some point, we immediately conclude. Otherwise, we assume w.l.o.g. that $u>0$ and we consider the principal eigenvalue $\lambda_{1}=\lambda_{1}\left(a^{+}\right)$of the problem

$$
\ddot{\varphi}+\lambda a^{+}(t) \varphi=0, \quad \varphi \in H_{0}^{1}(0, \tau)
$$

together with the corresponding positive eigenfunction, $\varphi_{1}=\varphi_{1}\left(a^{+}\right)$. Testing with $u$ the equation solved by $\varphi_{1}$ and with $\varphi_{1}$ the equation solved by $u$, after an integration one obtains

$$
\left[\dot{\varphi}_{1} u-\dot{u} \varphi_{1}\right]_{0}^{\tau}=\int_{0}^{\tau}\left(\ddot{\varphi}_{1} u-\varphi_{1} \ddot{u}\right)=\int_{0}^{\tau} a^{+} u \varphi_{1}\left(u^{2}-\lambda_{1}\right) .
$$

Since the left hand side, $\dot{\varphi}_{1}(\tau) u(\tau)-\dot{\varphi}_{1}(0) u(0)$, is negative (indeed, $\varphi_{1}$ is positive and vanishes at the end-points of the interval) we conclude that

$$
\min _{t \in[0, \tau]}|u(t)| \leq \sqrt{\lambda_{1}}
$$

as desired.
Our second result deals with the distribution of zeros of a solution on $[0, \tau]$. To state it precisely, recall the definitions of $c$ and $c_{\zeta}$ in (14) and fix $\zeta>0$ so small that

$$
\begin{equation*}
2\left\|a^{+}\right\|_{L^{\infty}}\left(c+c_{\zeta}\right) \zeta^{3}<1 \tag{25}
\end{equation*}
$$

(this is clearly possible, since $c_{\zeta}$ is bounded for $\zeta \rightarrow 0^{+}$). This will be the value of $\zeta$ used henceforth.

Lemma 3.2. Let $u$ be a nontrivial solution of (20) such that

$$
\begin{equation*}
\int_{0}^{\tau} \dot{u}^{2} \leq 2\left(c+c_{\zeta}\right) \tag{26}
\end{equation*}
$$

Then, $u$ has at most one zero in $[0, \zeta]$ and at most one zero in $[\tau-\zeta, \tau]$.

Proof. Assume, just to fix the ideas, that there are $t_{1}, t_{2} \in[0, \zeta]$, with $t_{1}<t_{2}$, such that $u\left(t_{1}\right)=u\left(t_{2}\right)=0$. Then, from (18), we have

$$
\|u\|_{L^{\infty}\left(t_{1}, t_{2}\right)}^{4} \leq \zeta^{2}\left(\int_{t_{1}}^{t_{2}} \dot{u}^{2}\right)^{2}
$$

Multiplying equation (20) by $u$ and integrating by parts on $\left[t_{1}, t_{2}\right]$, from the above estimate and (26), it follows that

$$
\begin{aligned}
\int_{t_{1}}^{t_{2}} \dot{u}^{2} & =\int_{t_{1}}^{t_{2}} a^{+} u^{4} \leq\left\|a^{+}\right\|_{L^{\infty}} \zeta^{3}\left(\int_{t_{1}}^{t_{2}} \dot{u}^{2}\right)^{2} \\
& \leq\left\|a^{+}\right\|_{L^{\infty}} \zeta^{3} \int_{0}^{\tau} \dot{u}^{2} \int_{t_{1}}^{t_{2}} \dot{u}^{2} \leq 2\left\|a^{+}\right\|_{L^{\infty}}\left(c+c_{\zeta}\right) \zeta^{3} \int_{t_{1}}^{t_{2}} \dot{u}^{2} .
\end{aligned}
$$

In view of the choice of $\zeta$ in (25), we obtain a contradiction (since $u \not \equiv 0$ ).
In the next results, we study the behavior of solutions on $[\tau, T]$. Our first lemma in this direction plays a crucial role in the rest of the paper, describing the behavior of solutions at the boundary of $[\tau, T]$ when the parameter $\mu$ is large.

Lemma 3.3. For every $\varepsilon>0$ and $R>0$, there exists $\widehat{\mu}>0$ such that, for any $\mu>\widehat{\mu}$ and for any solution $u$ of (21) such that $\|u\|_{L^{\infty}(\tau, T)} \leq K$ (see assumption (C3) at page 8),

$$
\left\{\begin{array} { l } 
{ \dot { u } ( \tau ^ { + } ) \geq - R } \\
{ u ( \tau ) \geq 0 }
\end{array} \quad \Longrightarrow \quad u ( \tau ) \leq \varepsilon \quad \text { and } \quad \left\{\begin{array}{l}
\dot{u}\left(\tau^{+}\right) \leq R \\
u(\tau) \leq 0
\end{array} \quad \Longrightarrow \quad-\varepsilon \leq u(\tau)\right.\right.
$$

(and analogous estimates hold for $u(T)$ ).
Proof. We will deal only with the case $\dot{u}\left(\tau^{+}\right) \geq-R, u(\tau) \geq 0$, the other being analogous. Let us fix $\varepsilon>0$ and assume, by contradiction, that $u(\tau)>\varepsilon$. A simple convexity argument implies that (provided $\varepsilon$ is small enough, so that $\varepsilon / R<T-\tau$ )

$$
u(t) \geq \varepsilon-R(t-\tau) \quad \text { for every } t \in[\tau, \tau+\varepsilon / R]
$$

and $u(t) \geq \varepsilon / 2$ for every $t \in[\tau, \bar{t}]$, with $\bar{t}=\tau+\varepsilon /(2 R)$. Hence, integrating twice equation (21), we have

$$
\begin{aligned}
u(\bar{t}) & \geq u(\tau)+\dot{u}\left(\tau^{+}\right)(\bar{t}-\tau)+\mu \frac{\varepsilon^{3}}{8} \int_{\tau}^{\bar{t}} \int_{\tau}^{t} a^{-}(s) d s d t \\
& \geq \frac{\varepsilon}{2}+\mu \frac{\varepsilon^{3}}{8} \int_{\tau}^{\bar{t}} \int_{\tau}^{t} a^{-}(s) d s d t
\end{aligned}
$$

which contradicts $\|u\|_{L^{\infty}(\tau-T, 0)} \leq K$ whenever $\mu$ is sufficiently large, since $\int_{\tau}^{t} a^{-}(s) d s>0$ by (7).

Our next result deals with the decay of the solutions, when $\mu \rightarrow+\infty$, on compact subintervals of $(\tau, T)$.

Lemma 3.4. For any $\delta \in(0,(T-\tau) / 2)$ there exists $C_{\delta}>0$ such that, for any $u$ solving (21), it holds that

$$
\begin{equation*}
|u(t)| \leq C_{\delta}\left(\frac{\max (|u(\tau)|,|u(T)|)}{\mu}\right)^{1 / 3}, \quad \text { for every } t \in[\tau+\delta, T-\delta] \tag{27}
\end{equation*}
$$

Proof. Let us fix $\delta \in(0,(T-\tau) / 2)$. By convexity arguments it holds

$$
|u(t)| \leq \max (|u(\tau+\delta)|,|u(T-\delta)|), \quad \text { for every } t \in[\tau+\delta, T-\delta]
$$

hence we need an estimate on $\max (|u(\tau+\delta)|,|u(T-\delta)|)$. Let us assume that

$$
\max (|u(\tau+\delta)|,|u(T-\delta)|)=|u(T-\delta)|
$$

the opposite case being the same. Under this hypothesis it holds

$$
u(T-\delta) \geq 0(\text { resp. } \leq 0) \quad \Longrightarrow \quad \dot{u}(T-\delta) \geq 0(\text { resp. } \leq 0)
$$

We will detail the proof in the case $u(T-\delta) \geq 0$; first of all we remark that $u(t) \geq u(T-\delta)$ on $[T-\delta, T]$. Integrating twice equation (21) we have

$$
\begin{aligned}
u(T) & \geq u(T-\delta)+\dot{u}(T-\delta) \delta+\mu u^{3}(T-\delta) \int_{T-\delta}^{T} \int_{T-\delta}^{t} a^{-}(s) d s d t \\
& \geq \mu u^{3}(T-\delta) \int_{T-\delta}^{T} \int_{T-\delta}^{t} a^{-}(s) d s d t
\end{aligned}
$$

Being the above integral strictly positive by (7), we take $C_{\delta}=\left(\int_{T-\delta}^{T} \int_{T-\delta}^{t} a^{-}(s) d s d t\right)^{-1 / 3}$.

### 3.2 Some local Nehari-type arguments

In this section, we collect some results for $H^{1}$ functions which are "almost" in $\mathcal{N}$ (recall the definition (13)) in the sense that $\left|\int_{0}^{\tau}\left(\dot{u}^{2}-a^{+} u^{4}\right)\right|$ and $|u(0)|,|u(\tau)|$ are small. We fix here

$$
\begin{equation*}
r=\left(32\left\|a^{+}\right\|_{L^{\infty}} \tau^{3}\right)^{-1 / 2} \tag{28}
\end{equation*}
$$

and this will be the value of $r$ used throughout the paper.
Our first lemma says that, for functions almost in $\mathcal{N}$, the value $\int_{0}^{\tau} \dot{u}^{2}$ is either very small or quite large. It is worth noticing that from its proof we can conclude that $r^{2}<4 c$ (compare with Remark 2.1).
Lemma 3.5. For every $\varepsilon>0$ small, there exists a constant $\delta_{\varepsilon}>0$ and with $\delta_{\varepsilon} \rightarrow 0$ for $\varepsilon \rightarrow 0^{+}$such that, for any $u \in H^{1}(0, \tau)$,

$$
\left\{\begin{array} { l } 
{ | \int _ { 0 } ^ { \tau } ( \dot { u } ^ { 2 } - a ^ { + } u ^ { 4 } ) | \leq \varepsilon } \\
{ | u ( 0 ) | , | u ( \tau ) | \leq \varepsilon }
\end{array} \quad \Longrightarrow \quad \left\{\begin{array}{l}
\int_{0}^{\tau} \dot{u}^{2} \leq \delta_{\varepsilon} \quad \text { if } \quad \int_{0}^{\tau} \dot{u}^{2} \leq r^{2} \\
\int_{0}^{\tau} \dot{u}^{2} \geq 2 r^{2} \quad \text { if } \quad \int_{0}^{\tau} \dot{u}^{2} \geq r^{2}
\end{array}\right.\right.
$$

Proof. Assume that, for some $\varepsilon>0$, the conditions on the left-hand side hold true. Then from (7), it holds

$$
\|u\|_{L^{\infty}(0, \tau)} \leq \varepsilon+\sqrt{\tau}\left(\int_{0}^{\tau} \dot{u}^{2}\right)^{1 / 2}
$$

hence, recalling the elementary inequality $(A+B)^{4} \leq 8\left(A^{4}+B^{4}\right), A, B \geq 0$,

$$
\int_{0}^{\tau} a^{+} u^{4} \leq 8\left\|a^{+}\right\|_{L^{\infty}}\left[\tau \varepsilon^{4}+\tau^{3}\left(\int_{0}^{\tau} \dot{u}^{2}\right)^{2}\right]
$$

Summing up,

$$
\int_{0}^{\tau} \dot{u}^{2} \leq C_{\varepsilon}+C\left(\int_{0}^{\tau} \dot{u}^{2}\right)^{2}
$$

where $C_{\varepsilon}=\varepsilon+8\left\|a^{+}\right\|_{L^{\infty}} \tau \varepsilon^{4}$ and $C=8\left\|a^{+}\right\|_{L^{\infty}} \tau^{3}$. Solving the second order inequality in $\int_{0}^{\tau} \dot{u}^{2}$ we obtain

$$
\text { either } \quad \int_{0}^{\tau} \dot{u}^{2} \leq \delta_{\varepsilon}^{-} \quad \text { or } \quad \int_{0}^{\tau} \dot{u}^{2} \geq \delta_{\varepsilon}^{+}
$$

where

$$
\delta_{\varepsilon}^{ \pm}=\left(1 \pm \sqrt{1-4 C C_{\varepsilon}}\right) /(2 C)
$$

Since, for $\varepsilon \rightarrow 0^{+}, \delta_{\varepsilon}^{-} \rightarrow 0$ and $\delta_{\varepsilon}^{+} \geq 1 /(2 C)=2 r^{2}$, the thesis follows.
Our second result concerns the presence of internal zeros for functions almost in $\mathcal{N}$.
Lemma 3.6. Let $\zeta$ be as in (25). Then there exist $\bar{\varepsilon}>0$ and $\omega>0$ such that, for any $u \in H^{1}(0, \tau)$,

$$
\left\{\begin{array}{l}
\left|\int_{0}^{\tau}\left(\dot{u}^{2}-a^{+} u^{4}\right)\right| \leq \bar{\varepsilon} \\
|u(0)|,|u(\tau)| \leq \bar{\varepsilon} \\
r^{2} \leq \int_{0}^{\tau} \dot{u}^{2} \leq 2\left(c+c_{\zeta}\right)
\end{array} \quad \Longrightarrow \quad|u(t)| \geq \omega \quad \text { for every } t \in[\zeta, \tau-\zeta]\right.
$$

Proof. Assume by contradiction that there exists a sequence $\left(u_{n}\right) \subset H^{1}(0, \tau)$ with

$$
\begin{equation*}
\left|\int_{0}^{\tau}\left(\dot{u}_{n}^{2}-a^{+} u_{n}^{4}\right)\right| \rightarrow 0, \quad\left|u_{n}(0)\right|,\left|u_{n}(\tau)\right| \rightarrow 0, \quad r^{2} \leq \int_{0}^{\tau} \dot{u}_{n}^{2} \leq 2\left(c+c_{\zeta}\right) \tag{29}
\end{equation*}
$$

and such that

$$
u_{n}\left(t_{n}\right) \rightarrow 0, \quad \text { for some } t_{n} \in[\zeta, \tau-\zeta]
$$

From the second and the third condition in (29), we infer (once more using (17)) that $u_{n}$ is bounded in $H^{1}(0, \tau)$, so that there exists $\bar{u} \in H^{1}(0, \tau)$ such that $u_{n} \rightarrow \bar{u}$ weakly in $H^{1}$ and uniformly. Moreover, $\bar{u} \in H_{0}^{1}(0, \tau), \bar{u}\left(t^{*}\right)=0$ for a suitable $t^{*} \in[\zeta, \tau-\zeta]$, and $\bar{u} \not \equiv 0$ (since from the first and the third condition in (29) we know that $\int_{0}^{\tau} a^{+} u_{n}^{4} \geq r^{2} / 2$, for $n$ large). Finally,

$$
\int_{0}^{\tau} \dot{\bar{u}}^{2} \leq \liminf _{n \rightarrow+\infty} \int_{0}^{\tau} \dot{u}_{n}^{2} \leq \liminf _{n \rightarrow+\infty} \int_{0}^{\tau} a^{+} u_{n}^{4}=\int_{0}^{\tau} a^{+} \bar{u}^{4}
$$

Set now

$$
\bar{\lambda}=\left(\frac{\int_{0}^{\tau} \dot{\bar{u}}^{2}}{\int_{0}^{\tau} a^{+} \bar{u}^{4}}\right)^{1 / 2} \leq 1
$$

then the function $\hat{u}=\bar{\lambda} \bar{u}$ lies in $\mathcal{Z}_{\zeta}$. Hence

$$
4 c_{\zeta} \leq \int_{0}^{\tau} \dot{\hat{u}}^{2} \leq \int_{0}^{\tau} \dot{\bar{u}}^{2} \leq \liminf _{n \rightarrow+\infty} \int_{0}^{\tau} \dot{u}_{n}^{2} \leq 2\left(c+c_{\zeta}\right)<4 c_{\zeta},
$$

a contradiction.

## 4 Some properties of $\mathcal{N}_{\mu}$ and proof of Proposition 2.4

In this section we establish some further fundamental properties enjoyed by functions in $\mathcal{N}_{\mu}$, as defined in equation (15), and, as a consequence, we give the proof of Proposition 2.4.

Contrarily to the properties in Lemma 2.3 , the fact that the parameter $\mu$ is large now plays a role. We point out that the final value $\mu^{*}$ in Theorem 2.2 will be the result of many successive enlargements; with some abuse of notation, but not to overload it, in all the subsequent results we will always use the same symbol $\mu^{*}$ to denote the outcome at each step. The crucial point, however, is that the rule for any of these enlargements depends ultimately only on the local estimate given in Lemma 3.3 (that is, on the value $\widehat{\mu}$ given therein). For this reason, the final value $\mu^{*}$ is a quantity independent on both the sequence $\mathcal{L}$ and on the integer $N$.

Lemma 4.1. For every $\varepsilon>0$ there exists $\mu^{*}>0$ such that, for any $\mu>\mu^{*}$ and $u \in \mathcal{N}_{\mu}$, the following hold true, for any $i \in\{-N, \ldots, N\}$ :
(i) $\left|u\left(\sigma_{i}\right)\right|+\left|u\left(\tau_{i}\right)\right| \leq \varepsilon$,
(ii) $\left|\dot{u}\left(\sigma_{i}^{-}\right) u\left(\sigma_{i}\right)\right|+\left|\dot{u}\left(\tau_{i}^{+}\right) u\left(\tau_{i}\right)\right| \leq \varepsilon$,
(iii) $\left|\int_{I_{i}^{+}}\left(\dot{u}^{2}-a^{+} u^{4}\right)\right| \leq \varepsilon$.

Proof. As for (i), when $i \in L$ assumption (C4) holds and the conclusion follows directly from Lemma 3.3, with the choice $R=\rho$. On the other hand, when $i \notin L$ we first apply Lemma 3.1 with $M=r^{2}$ to ensure that

$$
\begin{equation*}
\|\dot{u}\|_{L^{\infty}\left(I_{i}^{+}\right)} \leq M^{\prime} \tag{30}
\end{equation*}
$$

for a suitable $M^{\prime}>0$. Second, we notice that $\dot{u}\left(\sigma_{i}^{-}\right)=\dot{u}\left(\sigma_{i}^{+}\right)$and $\dot{u}\left(\tau_{i}^{+}\right)=\dot{u}\left(\tau_{i}^{-}\right)$since, by Lemma 2.3 (i), $u$ solves the differential equation both at $\sigma_{i}$ and $\tau_{i}$. Hence, the conclusion follows again from Lemma 3.3, with the choice $R=M^{\prime}$.

We now deal with (ii); again, we have to distinguish the cases $i \notin L$ and $i \in L$. In the former, the conclusion is immediate, recalling (30) (together with $\dot{u}\left(\sigma_{i}^{-}\right)=\dot{u}\left(\sigma_{i}^{+}\right)$and $\left.\dot{u}\left(\tau_{i}^{+}\right)=\dot{u}\left(\tau_{i}^{-}\right)\right)$and the previous point (i). In the latter, consider for instance the case $u\left(\tau_{i}\right)>0$ (the other being analogous). Since $u\left(\tau_{i}\right)$ is small by the previous step, and $\dot{u}\left(\tau_{i}^{+}\right)>$ $-\rho$, we just have to prove that $\dot{u}\left(\tau_{i}^{+}\right)$is bounded from above. Since $u$ is convex on $I_{i}^{-}$and $u$ satisfies (C3), we have

$$
\dot{u}\left(\tau_{i}^{+}\right) \leq \frac{u\left(\sigma_{i+1}\right)-u\left(\tau_{i}\right)}{\sigma_{i+1}-\tau_{i}} \leq \frac{K}{T-\tau}
$$

and the conclusion follows.
Finally, we notice that for every index $i$ it holds

$$
\int_{I_{i}^{+}}\left(\dot{u}^{2}-a^{+} u^{4}\right)=u\left(\tau_{i}\right) \dot{u}\left(\tau_{i}^{+}\right)-u\left(\sigma_{i}\right) \dot{u}\left(\sigma_{i}^{-}\right) .
$$

This equality comes from Lemma 2.3 (iii) if $i \in L$, while it is a consequence of the integration by parts rule when $i \notin L$. Then, (iii) follows from (ii).

Lemma 4.2. For every $\varepsilon>0$ small enough there exists $\mu^{*}>0$ such that, for any $\mu>\mu^{*}$ and $u \in \mathcal{N}_{\mu}$, the following hold true:
(i) $\|u\|_{L^{\infty}\left(I_{i}^{-}\right)}+\int_{I_{i}^{-}} \dot{u}^{2} \leq \varepsilon$, for any $i \in\{-N, \ldots, N\}$,
(ii) $\|u\|_{L^{\infty}\left(I_{i}^{+}\right)}+\|\dot{u}\|_{L^{\infty}\left(I_{i}^{+}\right)}+\|\ddot{u}\|_{L^{\infty}\left(I_{i}^{+}\right)} \leq \varepsilon$, for any $i \notin L$,
(iii) $\int_{I_{i}^{+}} \dot{u}^{2} \geq 2 r^{2}$, for any $i \in L$,
(iv) for a suitable $\omega>0$,

$$
u(t) \geq \omega, \quad \text { for every } t \in \bigcup_{i \in L}\left[\sigma_{i}+\zeta, \tau_{i}-\zeta\right]
$$

Proof. Integrating by parts equation (21) solved by $u$ on $I_{i}^{-}$, we have

$$
\begin{aligned}
\int_{I_{i}^{-}} \dot{u}^{2} & =-\mu \int_{I_{i}^{-}} a^{-} u^{4}+u\left(\sigma_{i+1}\right) \dot{u}\left(\sigma_{i+1}^{-}\right)-u\left(\tau_{i}\right) \dot{u}\left(\tau_{i}^{+}\right) \\
& \leq u\left(\sigma_{i+1}\right) \dot{u}\left(\sigma_{i+1}^{-}\right)-u\left(\tau_{i}\right) \dot{u}\left(\tau_{i}^{+}\right)
\end{aligned}
$$

Recalling (22), the conclusion (i) follows from the points (i) and (ii) of Lemma 4.1.
We now deal with (ii). The elementary inequality (17) together with Lemma 3.5 and the point (i) of Lemma 4.1 imply that $\|u\|_{L^{\infty}\left(I_{i}^{+}\right)}$can be made arbitrarily small for $\mu$ large. Since $u$ solves an equation independent of $\mu$ on such an interval, we deduce the bound for $\|\ddot{u}\|_{L^{\infty}\left(I_{i}^{+}\right)}$. Using (24), we conclude that the same is true for $\|\dot{u}\|_{L^{\infty}\left(I_{i}^{+}\right)}$, as well.

Finally, (iii) and (iv) follow from Lemmas 3.5 and 3.6, respectively.
We are now in position to prove Proposition 2.4. As we mentioned, such a proof relies on the main result of [24], which we report here for the reader's convenience.

Theorem 4.3 ([24, Theorems 1.2, 2.8]). Let $X$ be a Hilbert space, $J \in C^{2}(X, \mathbb{R}), V^{+} \subset X$ a fixed closed linear subspace. We define

$$
V_{x}^{+} \equiv V^{+}, \quad V_{x}^{-}=\operatorname{span}\left\{\xi_{1}(x), \ldots, \xi_{h}(x)\right\}, \quad V_{x}=V_{x}^{+} \oplus V_{x}^{-}
$$

with $\xi_{i} \in C^{1}(\mathcal{A}, X)$ for every $i=1, \ldots, h, \mathcal{A} \subset X$ open, in such a way that $V_{x}$ is a proper subspace for every x. Let

$$
\mathcal{M}=\left\{x \in \mathcal{A}: \operatorname{proj}_{V_{x}} \nabla J(x)=0\right\} .
$$

Let us suppose that for some $0<\delta^{\prime}<\delta^{\prime \prime}$ it holds, for every $x \in \mathcal{M}$,
(i) $\left\|\xi_{i}(x)\right\|_{X} \geq \delta^{\prime},\left\langle\xi_{i}(x), \xi_{j}(x)\right\rangle_{X}=0$, for every $i \neq j$;
(ii) $\xi_{i}^{\prime}(x)[v] \in V_{x}$ for every $i$ and $v \in V_{x}$;
(iii) $\pm J^{\prime \prime}(x)[v, v] \geq \delta^{\prime}\|v\|_{X}^{2}$ for every $v \in V_{x}^{ \pm}$;
(iv) $\left\|\xi_{i}^{\prime}(x)[u]\right\|_{X} \leq \delta^{\prime \prime}\|u\|_{X},\left|J^{\prime}(x)[u]\right| \leq \delta^{\prime \prime}\|u\|_{X}$ and $\left|J^{\prime \prime}(x)[u, w]\right| \leq \delta^{\prime \prime}\|u\|_{X}\|w\|_{X}$ for every $u, w \in X$.

Then the set $\mathcal{M}$ is a $C^{1}$ embedded submanifold of $X$ such that any constrained Palais-Smale sequence for $J$ is a free one.

Proof of Proposition 2.4. We argue in two steps. First of all, we consider the open subset of $H_{\text {per }}^{1}\left(I_{N}\right)$ defined as

$$
\mathcal{A}=\left\{\begin{array}{ll} 
& \int_{I_{i}^{+}} u^{2}<\delta^{2}, \text { for } i \notin L \\
u \in H_{\mathrm{per}}^{1}\left(I_{N}\right): & \int_{I_{i}^{+}} \dot{u}^{2}>r^{2}, \text { for } i \in L \\
& 3 \int_{I_{N}} \dot{\eta}_{i}^{2} u^{2}<r^{2}, \text { for } i \in L
\end{array}\right\}
$$

where

$$
\delta=\min \left\{r, \frac{1}{2}\left[3 T(k+1)\left\|a^{+}\right\|_{L^{\infty}}\right]^{-1 / 2}\right\}
$$

and, recalling definition (12), we set

$$
\mathcal{M}_{\mu}=\mathcal{A} \cap \mathcal{P}_{\mu}
$$

Then, we have the following.
Step 1. For any $\mu>0$, the set $\mathcal{M}_{\mu}$ is a $C^{1}$ embedded submanifold of $H_{\mathrm{per}}^{1}\left(I_{N}\right)$ such that any constrained Palais-Smale sequence is a free one.
Of course, this comes from a direct application of Theorem 4.3, with $X=H_{\mathrm{per}}^{1}\left(I_{N}\right), J=J_{\mu}$, $\mathcal{M}=\mathcal{M}_{\mu}$, and $\xi_{i}(u)=\eta_{i} u, i \in L$ (recall (12)). From this point of view, assumptions (i), (ii) and (iv) are almost straightforward, and we refer to the proof of [24, Theorem 3.5] for full details. Here we focus on the proof of (iii).

Let $u \in \mathcal{M}_{\mu}$ and $v \in V^{+}$; in particular, $v$ vanishes on $\cup_{i \in L} I_{i}^{+}$. We have

$$
J_{\mu}^{\prime \prime}(u)[v, v]=\int_{I_{N}} \dot{v}^{2}-3 \int_{I_{N}} a_{\mu} u^{2} v^{2} \geq \int_{I_{N}} \dot{v}^{2}-3 \int_{I_{N}} a^{+} u^{2} v^{2}
$$

We now indicate with $\Delta_{l}$ the connected components of $I_{N} \backslash \cup_{i \in L} I_{i}^{+}$(the index $l$ varies between 1 and the cardinality of $L,|L|$, or between 1 and $|L|+1$ ). Recalling assumption (8), the following estimate on the length of each $\Delta_{l}$ holds

$$
\left|\Delta_{l}\right| \leq k T+(T-\tau) \leq(k+1) T .
$$

Hence, since $v$ vanishes at the end-points of each $\Delta_{l}$, once more using (17), for every $k$ it holds

$$
\|v\|_{L^{\infty}\left(\Delta_{l}\right)}^{2} \leq\left|\Delta_{l}\right| \int_{\Delta_{l}} \dot{v}^{2} \leq(k+1) T \int_{\Delta_{l}} \dot{v}^{2},
$$

and hence

$$
\begin{aligned}
\int_{I_{N}} a^{+} u^{2} v^{2} & \leq \sum_{l} \int_{\Delta_{l}} a^{+} u^{2} v^{2} \leq \sum_{l}(k+1) T \int_{\Delta_{l}} \dot{v}^{2} \int_{\Delta_{l}} a^{+} u^{2} \\
& \leq(k+1) T \cdot\left\|a^{+}\right\|_{L^{\infty}} \cdot k \delta^{2} \sum_{l} \int_{\Delta_{l}} \dot{v}^{2} \\
& =k(k+1) T\left\|a^{+}\right\|_{L^{\infty}} \delta^{2} \int_{I_{N}} \dot{v}^{2}
\end{aligned}
$$

indeed at most $k$ intervals $I_{i}^{+}$with $i \notin L$ stay in each $\Delta_{l}$. Continuing the previous estimates on the second derivative of $J_{\mu}$ we obtain

$$
\begin{aligned}
J_{\mu}^{\prime \prime}(u)[v, v] & \geq\left(1-3 k(k+1) T\left\|a^{+}\right\|_{L^{\infty}} \delta^{2}\right) \int_{I_{N}} \dot{v}^{2} \\
& \geq\left(1+(k+1)^{2} T^{2}\right)^{-1}\left(1-3(k+1) T\left\|a^{+}\right\|_{L^{\infty}} k \delta^{2}\right)\|v\|_{H^{1}\left(I_{N}\right)}^{2} \\
& =\delta^{\prime}\|v\|_{H^{1}\left(I_{N}\right)}^{2}
\end{aligned}
$$

by the Poincaré inequality (19), where $\delta^{\prime}>0$ by the definition of $\delta$.

On the other hand, let $v \in V_{u}^{-}$. Then $v=\sum_{i \in L} \alpha_{i} \eta_{i} u$ for some $\alpha_{i} \in \mathbb{R}$ and, by the assumptions on $\eta_{i}, \dot{v}^{2}=\sum_{i \in L} \alpha_{i}^{2}\left(\dot{\eta_{i} u}\right)^{2}$. Using Lemma 2.3 (iv), we obtain

$$
\begin{aligned}
J_{\mu}^{\prime \prime}(u)[v, v] & =\int_{I_{N}} \dot{v}^{2}-3 \int_{I_{N}} a_{\mu}(t) u^{2} v^{2}=\sum_{i \in L} \alpha_{i}^{2}\left(\int _ { I _ { N } } \left(\dot{\left.\left.\eta_{i} u\right)^{2}-3 \int_{I_{N}} a_{\mu} \eta_{i}^{2} u^{4}\right)}\right.\right. \\
& =\sum_{i \in L} \alpha_{i}^{2}\left(-2 \int_{I_{N}}\left(\dot{\left.\eta_{i} u\right)^{2}}+3 \int_{I_{N}} \dot{\eta}_{i}^{2} u^{2}\right)\right. \\
& =-\int_{I_{N}} \dot{v}^{2}+\sum_{i \in L} \alpha_{i}^{2}\left(3 \int_{I_{N}} \dot{\eta}_{i}{ }^{2} u^{2}-\int_{I_{N}}\left(\dot{\eta}_{i} u\right)^{2}\right) .
\end{aligned}
$$

Since, for every $i \in L, u \in \mathcal{M}_{\mu}$,

$$
3 \int_{I_{N}}{\dot{\eta_{i}}}^{2} u^{2}<r^{2}<\int_{I_{i}^{+}} \dot{u}^{2}=\int_{I_{i}^{+}}\left(\dot{\left.\eta_{i} u\right)^{2}}<\int_{I_{N}}\left(\dot{\left.\eta_{i} u\right)^{2}}\right.\right.
$$

the conclusion follows again by Poincaré inequality.
Step 2. There exists $\mu^{*}>0$ such that, for $\mu>\mu^{*}, \mathcal{N}_{\mu} \subset \mathcal{M}_{\mu}$ open.
According to the definition of $\mathcal{A}$, we have to verify that, for $\mu$ large,

$$
\int_{I_{i}^{+}} u^{2}<\delta^{2}, \quad \text { for } i \notin L, \quad \text { and } \quad 3 \int_{I_{N}} \dot{\eta}_{i}^{2} u^{2}<r^{2}, \quad \text { for } i \in L
$$

the first inequality follows directly from the smallness of $\|u\|_{L^{\infty}\left(I_{i}^{+}\right)}$proved in Lemma 4.2 (ii). In the second one the integral is indeed on the support of $\dot{\eta}_{i}$ (which is contained in $I_{i-1}^{-} \cup I_{i}^{-}$), hence the thesis follows from the smallness of $\|u\|_{L^{\infty}\left(I_{i}^{-}\right)}$proved in Lemma 4.2 (i). Notice that the value $\mu^{*}$ obtained in this way depends (on the function $a$ and) on $k$, via $\delta$.

## 5 Construction of a local variation

In this section, we collect some results which will be used in the next Section 6 to construct local variations (in $\mathcal{N}_{\mu}$ ) for functions in $\partial \mathcal{N}_{\mu}$.

### 5.1 Variations on [0, $\tau]$

Let $r, \zeta$ be fixed as in (28) and (25). Take $p, q:[0, \tau] \rightarrow \mathbb{R}$ as the affine functions such that

$$
\begin{equation*}
q(\tau)=p(0)=0, \quad q(0)=p(\tau)=1 \tag{31}
\end{equation*}
$$

and define $\rho>0$ as

$$
\rho=16 \sup \left\{\frac{2 K}{T-\tau}+\left|\int_{0}^{\tau}\left(\dot{u} \dot{p}-a^{+} u^{3} p\right)\right|+\left|\int_{0}^{\tau}\left(\dot{u} \dot{q}-a^{+} u^{3} q\right)\right|: \begin{array}{l}
u \in H^{1}(0, \tau)  \tag{32}\\
|u(0)|,|u(\tau)| \leq 1 \\
\int_{0}^{\tau} \dot{u}^{2} \leq 2\left(c+c_{\zeta}\right)
\end{array}\right\}
$$

Remark 5.1. Choosing as a test function in the supremum above any element of $\mathcal{K}_{0}$ (recall definitions (4), (5)), an integration by parts yields $\left|\dot{u}\left(0^{+}\right)\right|+\left|\dot{u}\left(\tau^{-}\right)\right|<\rho$ for every $u \in \mathcal{K}_{0}$, and thus

$$
\left|\dot{u}\left(t^{ \pm}\right)\right| \leq \rho \quad \text { for every } u \in \mathcal{K}_{\mathcal{L}}, t \in \mathbb{R}
$$

The next result is quite technical but it greatly simplifies the exposition of some arguments in Section 6.

Lemma 5.1. There exists $\varepsilon>0$ such that, for any $\alpha_{0}^{\prime}, \beta_{0}^{\prime} \in[-\rho, \rho]$ and $u \in H^{1}(0, \tau)$ such that

$$
\left\{\begin{array}{l}
r^{2} \leq \int_{0}^{\tau} \dot{u}^{2} \leq 2\left(c+c_{\zeta}\right)  \tag{33}\\
|u(0)|,|u(\tau)| \leq \varepsilon \\
\int_{0}^{\tau}\left(\dot{u}^{2}-a^{+} u^{4}\right)=u(\tau) \beta_{0}^{\prime}-u(0) \alpha_{0}^{\prime}
\end{array}\right.
$$

there is a $C^{1}$-map defined on a neighborhood $O$ of $P_{0}=\left(u(0), u(\tau), \alpha_{0}^{\prime}, \beta_{0}^{\prime}\right)$

$$
U=U\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right): O \rightarrow H^{1}(0, \tau)
$$

such that:
(i) $U\left(P_{0}\right)(t)=u(t)$ for any $t \in[0, \tau]$ and, for every $\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right) \in O$,

$$
U\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right)(0)=\alpha, \quad U\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right)(\tau)=\beta
$$

(ii) for every $\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right) \in O$,

$$
\int_{0}^{\tau}\left(\dot{U}^{2}\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right)-a^{+} U^{4}\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right)\right)=\beta \beta^{\prime}-\alpha \alpha^{\prime}
$$

(iii) for a suitable $C>0$, depending only on $r, \zeta, \rho, \varepsilon$, it holds

$$
\begin{equation*}
\left\|\nabla_{\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right)} U\left(P_{0}\right)\right\|_{H^{1}(0, \tau)} \leq C \tag{34}
\end{equation*}
$$

Moreover, there exists $C_{\varepsilon}>0$, depending only on $r, \zeta, \rho, \varepsilon$ and such that $C_{\varepsilon} \rightarrow 0$ for $\varepsilon \rightarrow 0^{+}$, such that

$$
\begin{equation*}
\left|\partial_{\alpha} J_{[0, \tau]}\left(U\left(P_{0}\right)\right)\right| \leq \frac{\rho}{8}, \quad\left|\partial_{\beta} J_{[0, \tau]}\left(U\left(P_{0}\right)\right)\right| \leq \frac{\rho}{8} \tag{35}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\partial_{\alpha^{\prime}} J_{[0, \tau]}\left(U\left(P_{0}\right)\right)\right| \leq C_{\varepsilon}|u(0)|, \quad\left|\partial_{\beta^{\prime}} J_{[0, \tau]}\left(U\left(P_{0}\right)\right)\right| \leq C_{\varepsilon}|u(\tau)| \tag{36}
\end{equation*}
$$

(where $\left.J_{[0, \tau]}(u)=\int_{0}^{\tau} \frac{1}{2} \dot{u}^{2}-\frac{1}{4} a^{+} u^{4}\right)$.
Proof. We define

$$
F\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}, \lambda\right)=\int_{0}^{\tau}\left(\dot{V}^{2}(t ; \alpha, \beta, \lambda)-a^{+} V^{4}(t ; \alpha, \beta, \lambda)\right)-\left(\beta \beta^{\prime}-\alpha \alpha^{\prime}\right)
$$

where, for $t \in[0, \tau]$,

$$
V(t ; \alpha, \beta, \lambda)=\alpha q(t)+\beta p(t)+\lambda[u(t)-u(0) q(t)-u(\tau) p(t)] .
$$

By construction,

$$
V(t ; u(0), u(\tau), 1)=u(t), \quad V(0 ; \alpha, \beta, \lambda)=\alpha, \quad V(\tau ; \alpha, \beta, \lambda)=\beta
$$

moreover, $F\left(P_{0}, 1\right)=0$. Simple calculations show that

$$
\partial_{\lambda} F\left(P_{0}, 1\right)=\int_{0}^{\tau}\left(2 \dot{u}^{2}-4 a^{+} u^{4}\right)-\chi_{1}(u)
$$

where

$$
\chi_{1}(u)=u(0) \int_{0}^{\tau}\left(2 \dot{u} \dot{q}-4 a^{+} u^{3} q\right)+u(\tau) \int_{0}^{\tau}\left(2 \dot{u} \dot{p}-4 a^{+} u^{3} p\right) .
$$

Since

$$
\int_{0}^{\tau}\left(2 \dot{u}^{2}-4 a^{+} u^{4}\right)=-2 \int_{0}^{\tau} \dot{u}^{2}+4\left(u(\tau) \beta_{0}^{\prime}-u(0) \alpha_{0}^{\prime}\right),
$$

and, for $\varepsilon \rightarrow 0^{+}$,

$$
\left|u(\tau) \beta_{0}^{\prime}-u(0) \alpha_{0}^{\prime}\right|+\left|\chi_{1}(u)\right| \rightarrow 0
$$

uniformly in the class of functions satisfying (33), it holds that

$$
\partial_{\lambda} F\left(P_{0}, 1\right) \leq-2 r^{2}+o_{\varepsilon}(1) \leq-r^{2},
$$

whenever $\varepsilon>0$ is small enough (depending on $r, \zeta, \rho$, but not on $u$ ). Hence, the implicit function theorem applies, yielding the existence of $\lambda=\lambda\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right)$ solving implicitly $F\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}, \lambda\right)=0$ near $\left(P_{0}, 1\right)$. Then, with the position

$$
U\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right)(t)=V\left(t ; \alpha, \beta, \lambda\left(\alpha, \beta, \alpha^{\prime}, \beta^{\prime}\right)\right)
$$

the points (i) and (ii) of the statement are proved. We also notice that

$$
\partial_{\alpha} \lambda\left(P_{0}\right)=-\frac{\alpha_{0}^{\prime}+\int_{0}^{\tau}\left(2 \dot{u} \dot{q}-4 a^{+} u^{3} q\right)}{\partial_{\lambda} F\left(P_{0}, 1\right)} \quad \text { and } \quad \partial_{\alpha^{\prime}} \lambda\left(P_{0}\right)=-\frac{u(0)}{\partial_{\lambda} F\left(P_{0}, 1\right)} ;
$$

hence

$$
\begin{equation*}
\left|\partial_{\alpha} \lambda\left(P_{0}\right)\right| \leq C^{\prime} \quad \text { and } \quad\left|\partial_{\alpha^{\prime}} \lambda\left(P_{0}\right)\right| \leq C^{\prime}|u(0)|, \tag{37}
\end{equation*}
$$

where $C^{\prime}>0$ is a suitable constant depending only on $r, \zeta, \rho, \varepsilon$. Of course, similar estimates hold for the derivatives with respect to $\beta$ and $\beta^{\prime}$. From this consideration, (34) immediately follows.

To prove the final part of the statement, we observe that

$$
\begin{aligned}
\left.\frac{\partial}{\partial \alpha} J_{[0, \tau]}(U)\right|_{P_{0}} & =\int_{0}^{\tau}\left[\dot{U} \partial_{\alpha} \dot{U}-a^{+} U^{3} \partial_{\alpha} U\right]_{P_{0}} \\
& =\int_{0}^{\tau}\left[\dot{V}\left(\partial_{\alpha} \dot{V}+\partial_{\alpha} \lambda \partial_{\lambda} \dot{V}\right)-a^{+} V^{3}\left(\partial_{\alpha} V+\partial_{\alpha} \lambda \partial_{\lambda} V\right)\right]_{P_{0}} \\
& =\partial_{\alpha} \lambda\left(P_{0}\right) \int_{0}^{\tau}\left(\dot{V} \partial_{\lambda} \dot{V}-a^{+} V^{3} \partial_{\lambda} V\right)_{P_{0}}+\int_{0}^{\tau}\left(\dot{V} \partial_{\alpha} \dot{V}-a^{+} V^{3} \partial_{\alpha} V\right)_{P_{0}} \\
& =\partial_{\alpha} \lambda\left(P_{0}\right)\left[\int_{0}^{\tau}\left(\dot{u}^{2}-a^{+} u^{4}\right)-\chi_{2}(u)\right]+\int_{0}^{\tau}\left(\dot{u} \dot{q}-a^{+} u^{3} q\right),
\end{aligned}
$$

where

$$
\chi_{2}(u)=u(0) \int_{0}^{\tau}\left(\dot{u} \dot{q}-a^{+} u^{3} q\right)+u(\tau) \int_{0}^{\tau}\left(\dot{u} \dot{p}-a^{+} u^{3} p\right) .
$$

Since, for $\varepsilon \rightarrow 0^{+}$,

$$
\left|\int_{0}^{\tau}\left(\dot{u}^{2}-a^{+} u^{4}\right)\right|+\left|\chi_{2}(u)\right| \rightarrow 0
$$

uniformly in the class of functions satisfying (33), we have that (35) follows from the definition of $\rho$ in (32), and (37) (the estimate for the derivative w.r.t. $\beta$ can be obtained in the same way).

With even simpler computations we find

$$
\begin{aligned}
\left.\frac{\partial}{\partial \alpha^{\prime}} J_{[0, \tau]}(U)\right|_{P_{0}} & =\partial_{\alpha^{\prime}} \lambda\left(P_{0}\right) \int_{0}^{\tau}\left(\dot{V} \partial_{\lambda} \dot{V}-a^{+} V^{3} \partial_{\alpha} V_{\lambda}\right)_{P_{0}} \\
& =\partial_{\alpha^{\prime}} \lambda\left(P_{0}\right)\left[\int_{0}^{\tau}\left(\dot{u}^{2}-a^{+} u^{4}\right)-\chi_{2}(u)\right]
\end{aligned}
$$

so that, using similar arguments as above, (36) follows.

### 5.2 An auxiliary boundary value problem

In this section we establish some auxiliary results dealing with the following boundary value problem set on $[-T+\tau, T]=I_{-1}^{-} \cup I_{0}^{+} \cup I_{0}^{-}$

$$
\left\{\begin{array}{l}
\ddot{u}+a_{\mu}(t) u^{3}=0  \tag{38}\\
u(\tau-T)=x, u(T)=y \\
\int_{0}^{\tau} \dot{u}^{2}<r^{2}
\end{array}\right.
$$

where $r>0$ is as in (28). We postpone more comments about the role of the above written problem (38) in connection with the manifold $\mathcal{N}_{\mu}$ at the end of the Section (see Remark 5.2).

Proposition 5.2. For every $K>0$, there exists $\mu^{*}>0$ such that, for any $\mu>\mu^{*}$ and $x, y \in[-K, K]$, problem (38) admits a unique solution $\bar{u}_{\mu}(t ; x, y)$. Moreover,
(i) $\left\|\bar{u}_{\mu}\right\|_{L^{\infty}(\tau-T, T)} \leq K$,
(ii) uniformly in $x, y \in[-K, K]$,

$$
\begin{equation*}
\lim _{\mu \rightarrow+\infty}\left(\left\|\bar{u}_{\mu}\right\|_{L^{\infty}(0, \tau)}+\left\|\dot{\bar{u}}_{\mu}\right\|_{L^{\infty}(0, \tau)}\right)=0 \tag{39}
\end{equation*}
$$

(iii) if $\bar{u}_{\mu} \not \equiv 0$ and

$$
\begin{equation*}
\int_{\tau-T}^{T} \dot{\bar{u}}_{\mu}^{2}<\left((2 T-\tau)^{3}\left\|a^{+}\right\|_{L^{\infty}}\right)^{-1} \tag{40}
\end{equation*}
$$

then $\bar{u}_{\mu}$ cannot have neither more than one zero nor both a zero and a zero of the derivative in $[\tau-T, T]$. As a consequence: if $x y>0$, then $\bar{u}_{\mu}$ has constant sign on $[\tau-T, T]$; if $x y \leq 0$, then $\bar{u}_{\mu}$ vanishes exactly once in $[\tau-T, T]$ and $\dot{\bar{u}}_{\mu}(t) \neq 0$ for any $t \in[\tau-T, T]$.

Proof. We start by recalling the following: for every $K>0$ and $\varepsilon>0$, there exists $\mu^{*}>0$ such that, for any $\mu>\mu^{*}$ and $x, y \in[-K, K]$, then any solution $u$ of (38) satisfies

$$
\begin{equation*}
\|u\|_{L^{\infty}(0, \tau)}+\|\dot{u}\|_{L^{\infty}(0, \tau)} \leq \varepsilon \tag{41}
\end{equation*}
$$

This has already been proved in Lemma 4.2 (ii) and has many consequences. First of all, it implies (39); from this, point (i) follows using convexity arguments (see (22)). Second, it implies (iii). Indeed, if we assume that $t_{1}, t_{2} \in[\tau-T, T], t_{1}<t_{2}$, are instants satisfying one of the above properties, multiplying the equation $\ddot{u}+a_{\mu}(t) u^{3}=0$ by $u$ and integrating by parts on $\left[t_{1}, t_{2}\right]$ we find

$$
\int_{t_{1}}^{t_{2}} \dot{u}^{2}=\int_{t_{1}}^{t_{2}} a_{\mu} u^{4} \leq(2 T-\tau)\left\|a^{+}\right\|_{L^{\infty}}\|u\|_{L^{\infty}\left(t_{1}, t_{2}\right)}^{4}
$$

Using the Sobolev inequality (18) we obtain

$$
\int_{t_{1}}^{t_{2}} \dot{u}^{2} \leq(2 T-\tau)^{3}\left\|a^{+}\right\|_{L^{\infty}}\left(\int_{t_{1}}^{t_{2}} \dot{u}^{2}\right)^{2}
$$

contradicting (40).
Finally, also the uniqueness of the solution to (38) follows from (41). Indeed, assume that $u_{i}, i=1,2$, are two different solutions of (38) and define $w=u_{2}-u_{1}$; of course $w$ satisfies

$$
\left\{\begin{array}{l}
\ddot{w}+a_{\mu}(t)\left(u_{2}^{2}+u_{1} u_{2}+u_{1}^{2}\right) w=0 \\
w(\tau-T)=w(T)=0
\end{array}\right.
$$

Multiplying the previous equation by $w^{-}$, after an integration on $[\tau-T, T]$ we obtain

$$
\begin{aligned}
\int_{\tau-T}^{T}\left(\dot{w}^{-}\right)^{2} & =\int_{\tau-T}^{T} a_{\mu}\left(u_{2}^{2}+u_{1} u_{2}+u_{1}^{2}\right)\left(w^{-}\right)^{2} \\
& \leq \int_{[0, \tau] \cap \operatorname{supp}\left(w^{-}\right)} a^{+}\left(u_{2}^{2}+u_{1} u_{2}+u_{1}^{2}\right)\left(w^{-}\right)^{2}
\end{aligned}
$$

from (41) and the Poincaré inequality (19)

$$
\int_{\tau-T}^{T}\left(\dot{w}^{-}\right)^{2} \leq \varepsilon^{2}(2 T-\tau)^{2}\left\|a^{+}\right\|_{L^{\infty}} \int_{\tau-T}^{T}\left(\dot{w}^{-}\right)^{2}
$$

As $\mu$ is sufficiently large we obtain that necessarily $w^{-} \equiv 0$, that is $u_{2}-u_{1} \geq 0$. Exchanging the roles of $u_{1}$ and $u_{2}$, we infer $u_{1} \equiv u_{2}$.

We now turn to the existence, which is obtained via a variational argument and is organized in three steps.
Step 1. Reduction to a constrained minimization problem.
Consider the minimization problem

$$
\begin{equation*}
\inf _{u \in \mathcal{M}(x, y)} \widehat{J}_{\mu}(u) \tag{42}
\end{equation*}
$$

where

$$
\widehat{J}_{\mu}(u)=\frac{1}{2} \int_{\tau-T}^{T} \dot{u}^{2}-\frac{1}{4} \int_{\tau-T}^{T} a_{\mu} u^{4}, \quad u \in H^{1}(\tau-T, T)
$$

and

$$
\mathcal{M}(x, y)=\left\{\begin{array}{l|l}
u(\tau-T)=x, u(T)=y \\
u \in H^{1}(\tau-T, T) \mid & |u(0)| \leq K \\
& \int_{0}^{\tau} \dot{u}^{2} \leq r^{2}
\end{array}\right\}
$$

We observe that, for any $\mu>0$, such a minimization problem is for sure solvable. Indeed, the constraints $|u(0)| \leq K$ and $\int_{0}^{\tau} \dot{u}^{2} \leq r^{2}$ imply that

$$
\begin{equation*}
\|u\|_{L^{\infty}(0, \tau)} \leq K+\sqrt{\tau} r, \quad \text { for every } u \in \mathcal{M}(x, y) \tag{43}
\end{equation*}
$$

Hence, $\widehat{J}_{\mu}$ is coercive (and weakly lower semicontinuous) on the closed convex set $\mathcal{M}(x, y)$ and the conclusion follows from classical arguments.

Of course (by writing the Euler-Lagrange equation associated with (42)) we can prove that a solution of (42) solves (38) whenever $|u(0)|<K$ and $\int_{0}^{\tau} \dot{u}^{2}<r^{2}$. Steps 2 and 3 below will then be devoted to prove that this is the case when $\mu$ is large enough. In the following, we use the notation

$$
K_{\left[t_{1}, t_{2}\right]}(u)=\frac{1}{2} \int_{t_{1}}^{t_{2}} \dot{u}^{2}, \quad U_{\mu,\left[t_{1}, t_{2}\right]}(u)=\frac{1}{4} \int_{t_{1}}^{t_{2}} a_{\mu} u^{4}
$$

and

$$
\widehat{J}_{\mu,\left[t_{1}, t_{2}\right]}(u)=K_{\left[t_{1}, t_{2}\right]}(u)-U_{\mu,\left[t_{1}, t_{2}\right]}(u)
$$

for a subinterval $\left[t_{1}, t_{2}\right] \subset[\tau-T, T]$.
As a preliminary observation for the arguments below, we also notice that every minimizer of (42) solves the equation $\ddot{u}-\mu a^{-}(t) u^{3}=0$ on the interval $[\tau-T, 0] \cup[\tau, T]$. Indeed, variations vanishing on $[0, \tau]$ are admissible for functions in $\mathcal{M}(x, y)$.
Step 2. For $\mu$ large, the minimizers of (42) satisfy $|u(0)|<K$.
Let $u=u_{\mu}$ be a minimizer of (42) and assume by contradiction that $|u(0)|=K$. Our aim is to construct a function $\widetilde{u} \in \mathcal{M}(x, y)$ such that

$$
\begin{equation*}
\widehat{J}_{\mu}(\widetilde{u})<\widehat{J}_{\mu}(u), \tag{44}
\end{equation*}
$$

and thus contradicting the minimality of $u$.
To this end, we first construct $\widetilde{u}$ on $[\tau-T, 0]$. Let $\delta>0$ be small; according to Lemma 3.4, we have that $|u(t)|<K / 4$ for every $t \in[\tau-T+\delta,-\delta]$, provided $\mu$ is chosen so large that

$$
\begin{equation*}
C_{\delta}\left(\frac{K}{\mu}\right)^{1 / 3}<\frac{K}{4} \tag{45}
\end{equation*}
$$

with $C_{\delta}$ the constant appearing in (27). Then, we choose three points $t_{1}, t_{2}$ and $t^{*}$ in $[\tau-T, 0]$ as follows: $t_{1}$ is the minimum of $|u|$ (hence $\left.\left|u\left(t_{1}\right)\right|<K / 4\right) ; t_{2} \geq t_{1}$ is the unique point such that

$$
u\left(t_{2}\right)=2 u\left(t_{1}\right)
$$

(hence $\left.\left|u\left(t_{2}\right)\right|<K / 2\right) ; t^{*}>t_{2}$ is such that

$$
\left|u\left(t^{*}\right)\right|=\frac{K}{2}
$$

Notice that $t^{*} \geq-\delta$. Set now

$$
\widetilde{u}(t)= \begin{cases}u(t) & \text { if } t \in\left[\tau-T, t_{1}\right] \\ 2 u\left(t_{1}\right)-u(t) & \text { if } t \in\left[t_{1}, t_{2}\right] \\ 0 & \text { if } t \in\left[t_{2}, 0\right]\end{cases}
$$

(of course, the interval $\left[\tau-T, t_{1}\right]$ is empty if the minimum of $|u|$ is achieved at $\tau-T$, while the interval $\left[t_{1}, t_{2}\right]$ is empty if $u$ changes sign on $[\tau-T, 0]$, namely $\left.u\left(t_{1}\right)=0\right)$. Notice that:

- on $\left[\tau-T, t_{1}\right], u=\widetilde{u}$, so that $\widehat{J}_{\mu,\left[\tau-T, t_{1}\right]}(u)=\widehat{J}_{\mu,\left[\tau-T, t_{1}\right]}(\widetilde{u})$;
- on $\left[t_{1}, t_{2}\right],|\dot{u}|=|\dot{\widetilde{u}}|$ and $|u| \geq|\widetilde{u}|$, so that $K_{\left[t_{1}, t_{2}\right]}(u)=K_{\left[t_{1}, t_{2}\right]}(\widetilde{u})$ and $U_{\mu,\left[t_{1}, t_{2}\right]}(u) \leq$ $U_{\mu,\left[t_{1}, t_{2}\right]}(\widetilde{u})$; hence $\widehat{J}_{\mu,\left[t_{1}, t_{2}\right]}(u) \geq \widehat{J}_{\mu,\left[t_{1}, t_{2}\right]}(\widetilde{u})$
- on $\left[t_{2}, 0\right], \widehat{J}_{\mu,\left[t_{2}, 0\right]}(\widetilde{u})=0$ while (since $\left.K / 2=|u(0)|-\left|u\left(t^{*}\right)\right| \geq \int_{t^{*}}^{0}|\dot{u}|\right)$

$$
\begin{equation*}
\widehat{J}_{\mu,\left[t_{2}, 0\right]}(u) \geq \frac{1}{2} \int_{t^{*}}^{0} \dot{u}^{2} \geq \frac{K^{2}}{8}\left|t^{*}\right|^{-1} \geq \frac{K^{2}}{8} \delta^{-1} \tag{46}
\end{equation*}
$$

To construct $\widetilde{u}$ on $[0, T]$, we distinguish two possibilities. Set

$$
K^{\prime}=r \sqrt{\tau}
$$

If $|u(\tau)|<K^{\prime}$, we define

$$
\widetilde{u}(t)= \begin{cases}u(\tau) \frac{t}{\tau} & \text { if } t \in[0, \tau] \\ u(t) & \text { if } t \in[\tau, T] .\end{cases}
$$

Notice that

$$
\int_{0}^{\tau} \dot{\widetilde{u}}^{2}<\frac{\left(K^{\prime}\right)^{2}}{\tau}=r^{2}
$$

so that $\widetilde{u} \in \mathcal{M}(x, y)$. If $|u(\tau)| \geq K^{\prime}$, we set $\widetilde{u}(t)=0$ for $t \in[0, \tau]$ and then define $\widetilde{u}(t)$ on $[\tau, T]$ in a similar way to what has been done on $[\tau-T, 0]$ (with $K$ replaced by $K^{\prime}$ ). In any case, we have that:

- $\widehat{J}_{\mu,[\tau, T]}(u) \geq \widehat{J}_{\mu,[\tau, T]}(\widetilde{u})$;
- in view of (43), both $\widehat{J}_{\mu,[0, \tau]}(u)$ and $\widehat{J}_{\mu,[0, \tau]}(\widetilde{u})$ are bounded below by a constant independent of $u$ and $\mu$.

At this point, it is enough to observe that for $\delta \rightarrow 0^{+}$the term in (46) goes to infinity. In view of (45), we thus have a contradiction for $\mu$ large enough.
Step 3. For $\mu$ large, the minimizers of (42) satisfy $\int_{0}^{\tau} \dot{u}^{2}<r^{2}$.
Before starting the proof, we observe the following: for every $\varepsilon>0$, if $\mu$ is large enough any minimizer of (42) satisfies $|u(0)|,|u(\tau)| \leq \varepsilon$. This can be shown using the very same arguments of the previous Step 2. Hence, arguing as in the proof of Lemma 3.5 - and recalling the choice of $r$ in (28) - of we can see that the following relation holds

$$
\begin{equation*}
\int_{0}^{\tau} a^{+} u^{4} \leq K_{\varepsilon}+\frac{1}{4 r^{2}}\left(\int_{0}^{\tau} \dot{u}^{2}\right)^{2} \tag{47}
\end{equation*}
$$

where $K_{\varepsilon}=8\left\|a^{+}\right\|_{L^{\infty}} \tau \varepsilon^{4}$.
Now, let $u=u_{\mu}$ be a minimizer of (42) and assume by contradiction that $\int_{0}^{\tau} \dot{u}^{2}=r^{2}$; as in Step 2, we aim at constructing $\widetilde{u} \in \mathcal{M}(x, y)$ satisfying (44). Here we simply define $\widetilde{u}(t)=u(t)$ for $t \notin[0, \tau]$ and $\left.\widetilde{u}\right|_{[0, \tau]}$ as the affine functions for $(0, u(0))$ and $(\tau, u(\tau))$. Since $|u(0)|,|u(\tau)| \leq \varepsilon$, we see that $\widetilde{u} \in \mathcal{M}(x, y)$ for $\varepsilon$ small enough and

$$
\begin{equation*}
\widehat{J}_{\mu,[0, \tau]}(\widetilde{u})=o(1), \quad \text { for } \varepsilon \rightarrow 0^{+} . \tag{48}
\end{equation*}
$$

On the other hand, taking into account (47) we have

$$
\begin{equation*}
\widehat{J}_{\mu,[0, \tau]}(u) \geq \frac{r^{2}}{2}\left(1-\frac{1}{8}\right)-\frac{K_{\varepsilon}}{4}=\frac{7}{14} r^{2}-\frac{K_{\varepsilon}}{4}>0 \tag{49}
\end{equation*}
$$

for $\varepsilon$ small enough. Combining (48) and (49), we have that (44) holds true for $\mu$ large enough, as desired.

In the next proposition we collect some useful properties of the solution $\bar{u}_{\mu}(t ; x, y)$.
Proposition 5.3. Under the assumptions of Proposition 5.2, the map $(x, y) \mapsto \bar{u}_{\mu}(t ; x, y)$ $\in W^{2, \infty}(\tau-T, T)$ is of class $C^{1}$; setting

$$
v_{\mu}(t ; x, y)=\frac{\partial}{\partial x} \bar{u}_{\mu}(t ; x, y), \quad z_{\mu}(t ; x, y)=\frac{\partial}{\partial y} \bar{u}_{\mu}(t ; x, y)
$$

the following hold true (up to enlarging $\mu^{*}$ if necessary):
(i) $v_{\mu}, z_{\mu}$ solve, respectively, the linear boundary value problems

$$
\left\{\begin{array} { l } 
{ \ddot { v } + 3 a _ { \mu } ( t ) \overline { u } _ { \mu } ^ { 2 } ( t ; x , y ) v = 0 }  \tag{50}\\
{ v ( \tau - T ) = 1 , v ( T ) = 0 }
\end{array} \quad \left\{\begin{array}{l}
\ddot{z}+3 a_{\mu}(t) \bar{u}_{\mu}^{2}(t ; x, y) z=0 \\
z(\tau-T)=0, z(T)=1
\end{array}\right.\right.
$$

and, moreover, $v_{\mu}$ is positive and decreasing and $z_{\mu}$ is positive and increasing;
(ii) it holds

$$
\begin{equation*}
\frac{\partial}{\partial x} J_{\mu}\left(\bar{u}_{\mu}(\cdot ; x, y)\right)=-\dot{\bar{u}}_{\mu}\left((\tau-T)^{+} ; x, y\right) \tag{51}
\end{equation*}
$$

and

$$
\frac{\partial}{\partial y} J_{\mu}\left(\bar{u}_{\mu}(\cdot ; x, y)\right)=\dot{\bar{u}}_{\mu}\left(T^{-} ; x, y\right) ;
$$

(iii) it holds

$$
\begin{equation*}
\left|\dot{v}_{\mu}\left(T^{-} ; x, y\right)\right|,\left|\dot{z}_{\mu}\left((\tau-T)^{+} ; x, y\right)\right| \leq \frac{2}{2 T-\tau} \tag{52}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|x \dot{v}_{\mu}\left((\tau-T)^{+} ; x, y\right)\right|,\left|y \dot{z}_{\mu}\left(T^{-} ; x, y\right)\right| \leq \frac{2 K}{2 T-\tau}+5\left\|\dot{\bar{u}}_{\mu}\right\|_{L^{\infty}(T-\tau, T)} \tag{53}
\end{equation*}
$$

moreover, if $\max \left(\left|\dot{\bar{u}}_{\mu}\left((\tau-T)^{+} ; x, y\right)\right|,\left|\dot{\bar{u}}_{\mu}\left(T^{-} ; x, y\right)\right|\right) \leq \rho$, then

$$
\begin{equation*}
\left|x \dot{v}_{\mu}\left((\tau-T)^{+} ; x, y\right)\right|,\left|y \dot{z}_{\mu}\left(T^{-} ; x, y\right)\right| \leq \frac{2 K}{2 T-\tau}+5 \rho \tag{54}
\end{equation*}
$$

(iv) if $\bar{u}_{\mu}$ has constant sign on $[\tau-T, T]$, and

$$
\begin{array}{ll}
-\dot{\bar{u}}_{\mu}\left((\tau-T)^{+} ; x, y\right)=\dot{\bar{u}}_{\mu}\left(T^{-} ; x, y\right)=\rho & \text { if } \bar{u}_{\mu}>0 \\
\dot{\bar{u}}_{\mu}\left((\tau-T)^{+} ; x, y\right)=-\dot{\bar{u}}_{\mu}\left(T^{-} ; x, y\right)=\rho & \text { if } \bar{u}_{\mu}<0
\end{array}
$$

then it holds

$$
\begin{equation*}
\dot{v}_{\mu}\left((\tau-T)^{+} ; x, y\right)+\dot{z}_{\mu}\left((\tau-T)^{+} ; x, y\right)<0<\dot{v}_{\mu}\left(T^{-} ; x, y\right)+\dot{z}_{\mu}\left(T^{-} ; x, y\right) \tag{55}
\end{equation*}
$$

In points (iii) and (iv) of the above statement, $\rho$ is as in (32).
Proof. As already proved, $\bar{u}_{\mu}$ is unique; in the same spirit, one can use property (41), which says that the minimization problem is set in an almost convex case, to show that $\bar{u}_{\mu}$ is also non degenerate (i.e. that the linearized equation, with homogeneous Dirichlet boundary conditions, has no nontrivial solution). Using this and the Fredholm's Alternative, the $C^{1}-$ dependence from $x$ and $y$ follows in a standard way (see [25, Section 5] for the full details in a similar situation).

Now we prove separately each point of the statement; in points (i),(ii),(iii) we concentrate our attention on $v_{\mu}$, the proof for $z_{\mu}$ being analogous (for simplicity of notation, we omit the dependence on $x$ and $y$ if no confusion is possible).

The fact that $v_{\mu}$ solves the boundary value problem in (50) is well known; we prove that $v_{\mu}$ is positive and decreasing. Multiplying the equation in (50) by $v_{\mu}$ and integrating by parts, we find for any $t \in[\tau-T, T[$

$$
v_{\mu}(t) \dot{v}_{\mu}(t)=-\int_{t}^{T} \dot{v}_{\mu}^{2}+3 \int_{t}^{T} a_{\mu} \bar{u}_{\mu}^{2} v^{2} \leq-\int_{t}^{T} \dot{v}_{\mu}^{2}+3\left\|a^{+}\right\|_{L^{\infty}}\left\|\bar{u}_{\mu}\right\|_{L^{\infty}(0, \tau)}^{2} \int_{t}^{T} v_{\mu}^{2}
$$

By (19) we have

$$
v_{\mu}(t) \dot{v}_{\mu}(t) \leq\left(-1+3(2 T-\tau)^{2}\left\|a^{+}\right\|_{L^{\infty}}\left\|\bar{u}_{\mu}\right\|_{L^{\infty}(0, \tau)}^{2}\right) \int_{t}^{T} \dot{v}_{\mu}^{2}
$$

Recalling (39), we finally obtain, for $\mu$ large,

$$
v_{\mu}(t) \dot{v}_{\mu}(t) \leq-\frac{1}{2} \int_{t}^{T} \dot{v}_{\mu}^{2}<0, \quad \text { for every } t \in[\tau-T, T]
$$

and this proves our claim.
To prove (51), we simply observe that - integrating by parts -

$$
\frac{\partial}{\partial x} J_{\mu}\left(\bar{u}_{\mu}\right)=\int_{\tau-T}^{T}\left(\dot{\bar{u}}_{\mu} \dot{v}_{\mu}-a_{\mu} \bar{u}_{\mu}^{3} v_{\mu}\right)=\left[\dot{\bar{u}}_{\mu} v_{\mu}\right]_{\tau-T}^{T}-\int_{\tau-T}^{T}\left(\ddot{\bar{u}}_{\mu}+a_{\mu} \bar{u}_{\mu}^{3}\right) v_{\mu},
$$

whence the conclusion.
As for (52), we first use the mean value theorem to find $\bar{t} \in[\tau-T, T]$ such that $\dot{v}_{\mu}(\bar{t})=$ $-1 /(2 T-\tau)$. From this, and using the fact that $0 \leq v_{\mu} \leq 1$,

$$
\begin{aligned}
\dot{v}_{\mu}(T) & =\dot{v}_{\mu}(\bar{t})+\int_{\bar{t}}^{T} \ddot{v}_{\mu}=-\frac{1}{2 T-\tau}-3 \int_{\bar{t}}^{T} a_{\mu} \bar{u}_{\mu}^{2} v_{\mu} \\
& \geq-\frac{1}{2 T-\tau}-3\left\|a^{+}\right\|_{L^{\infty}}\left\|\bar{u}_{\mu}\right\|_{L^{\infty}(0, \tau)}^{2} .
\end{aligned}
$$

Recalling (39), we conclude for $\mu$ large. As for (53), taking into account the equations for $\bar{u}_{\mu}$ and $v_{\mu}$, we obtain

$$
3 \ddot{\bar{u}}_{\mu}(t) v_{\mu}(t)-\bar{u}_{\mu}(t) \ddot{v}_{\mu}(t)=0 .
$$

Integrating by parts we have

$$
-3 \dot{\bar{u}}_{\mu}(\tau-T)+x \dot{v}_{\mu}(\tau-T)-y \dot{v}_{\mu}(T)-2 \int_{\tau-T}^{T} \dot{\bar{u}}_{\mu} \dot{v}_{\mu}=0
$$

Now, the term $y \dot{v}_{\mu}(T)$ can be estimated using (52), while for the integral we have

$$
\left|\int_{\tau-T}^{T} \dot{\bar{u}}_{\mu}(t) \dot{v}_{\mu}(t)\right| \leq-\left\|\dot{\bar{u}}_{\mu}\right\|_{L^{\infty}(\tau-T, T)} \int_{\tau-T}^{T} \dot{v}_{\mu}=\left\|\dot{\bar{u}}_{\mu}\right\|_{L^{\infty}(\tau-T, T)} .
$$

This proves (53). As for (54), we just need to observe that

$$
\max \left(\left|\dot{\bar{u}}_{\mu}(\tau-T ; x, y)\right|,\left|\dot{\bar{u}}_{\mu}(T ; x, y)\right|\right) \leq \rho \quad \Longrightarrow \quad\left\|\dot{\bar{u}}_{\mu}\right\|_{L^{\infty}(\tau-T, T)} \leq \rho .
$$

This follows from (39), together with convexity arguments (see (23)).

Finally, we prove (55); we concentrate on the case $\bar{u}_{\mu}>0$, the other being analogous. Set $w_{\mu}=v_{\mu}+z_{\mu}$; then $w_{\mu}$ solves the boundary value problem

$$
\left\{\begin{array}{l}
\ddot{w}+3 a_{\mu}(t) \bar{u}_{\mu}^{2}(t ; x, y) w=0 \\
w(\tau-T)=w(T)=1
\end{array}\right.
$$

it holds $0 \leq w_{\mu} \leq 2$ and our thesis reads as $\dot{w}_{\mu}(\tau-T)<0<\dot{w}_{\mu}(T)$. Let us assume by contradiction that $\dot{w}_{\mu}(\tau-T) \geq 0$ (the argument for $\dot{w}_{\mu}(T)$ being the same); then, usual convexity arguments imply that $w_{\mu}(t) \geq 1$ for $t \in[\tau-T, 0]$ and that there exists $t^{*} \in[0, \tau]$ such that

$$
\dot{w}_{\mu}\left(t^{*}\right)=0
$$

(indeed, $\dot{w}_{\mu}(0)>0$ and $w_{\mu}(0)>1$; hence, if $\dot{w}_{\mu}(t)>0$ for every $t \in[0, \tau]$, then also $w_{\mu}(\tau)>1$ and, by convexity, $w_{\mu}(T)>1$, a contradiction).

Now, by Proposition 5.2, (i), $\bar{u}_{\mu} \leq K$ on $[\tau-T, 0]$; being $w_{\mu} \geq 1$ on the same interval, we deduce that $3 \ddot{\bar{u}}_{\mu} \leq K \ddot{w}_{\mu}$. Hence, for every $t \in[\tau-T, 0]$,

$$
\begin{aligned}
\dot{\bar{u}}_{\mu}(t) & =\dot{\bar{u}}_{\mu}(\tau-T)+\int_{\tau-T}^{t} \ddot{\bar{u}}_{\mu} \leq-\rho+\frac{K}{3} \int_{\tau-T}^{t} \ddot{w}_{\mu} \\
& =-\rho+\frac{K}{3}\left(\dot{w}_{\mu}(t)-\dot{w}_{\mu}(\tau-T)\right) \leq-\rho+\frac{K}{3} \dot{w}_{\mu}(t)=-\rho-\frac{K}{3} \int_{t}^{t^{*}} \ddot{w}_{\mu} \\
& \leq-\rho+K \int_{t}^{t^{*}} a_{\mu} \bar{u}_{\mu}^{2} w_{\mu} \leq-\rho+2 K\left\|a^{+}\right\|_{L^{\infty}}\left\|\bar{u}_{\mu}\right\|_{L^{\infty}(0, \tau)}^{2}
\end{aligned}
$$

In particular, according to (39) we can assume that $\dot{\bar{u}}_{\mu}(t) \leq-\rho / 2$ for every $t \in[\tau-T, 0]$ provided $\mu$ is large enough. Hence (recall (32))

$$
\bar{u}_{\mu}(0) \leq K-\frac{\rho}{2}(T-\tau) \leq 0
$$

contradicting the fact that $\bar{u}_{\mu}>0$.
We conclude this section with an important remark.
Remark 5.2. We observe at first that an existence and uniqueness result holds true, for any $\mu>0$, for the boundary value problem

$$
\left\{\begin{array}{l}
\ddot{u}-\mu a^{-}(t) u^{3}=0  \tag{56}\\
u(\tau)=x, u(T)=y
\end{array}\right.
$$

moreover (by convexity arguments), nontrivial solutions to (56) are of constant sign if $x y>0$, and are strictly monotone and vanish exactly once if $x y \leq 0$. The proof of these facts is straightforward, due to the coercivity and convexity of the action functional associated with (56). Moreover, a result completely analogous to Proposition 5.3 can be proved for (56) with simpler arguments (using, again, the convexity properties of the solutions). From this point view, Propositions 5.2 and 5.3 can be viewed as an extension of the more elementary results holding for (56) to the boundary value problem (38). Roughly speaking, it can be said that - as the integral bound $\int_{0}^{\tau} \dot{u}^{2}<r^{2}$ is considered - problem (38) behaves exactly like (56) when $\mu$ is large (notice, indeed, that - by point (ii) of Proposition 5.2 - solutions to (38) are arbitrarily small on $[0, \tau]$ when $\mu \rightarrow+\infty)$.

We now claim that Propositions 5.2 and 5.3 can be generalized for the problem

$$
\left\{\begin{array}{l}
\ddot{u}+a_{\mu}(t) u^{3}=0  \tag{57}\\
u\left(\tau_{i}\right)=x, u\left(\sigma_{i+l+1}\right)=y \\
\int_{\sigma_{j}}^{\tau_{j}} \dot{u}^{2}<r^{2} \quad \text { for } j=i+1, \ldots, i+l
\end{array}\right.
$$

(of course, (38) corresponds to the case $i=-1$ and $l=1$ ). The proofs can be obtained following exactly the same arguments given before, at the only expense of an unpleasant overloading of the notation. An important warning, however, is that in such a case the value $\mu^{*}$, as well as the constants in (52), (53) and (54), are depending on the integer $l$, that is, on the number of intervals of positivity of the weight function $a_{\mu}$ in the interval $\left[\tau_{i}, \sigma_{i+l+1}\right]$. Accordingly, all these constants can me made uniform for all problems of the type (57) with $l$ less than or equal to a common bound $k$ (compare with assumption (8)).

It now should be clear that Proposition 5.2 can be viewed as a tool for characterizing (and constructing) functions in $\mathcal{N}_{\mu}$ on intervals of the type

$$
I_{i}^{-} \cup I_{i+1}^{+} \cup I_{i+1}^{-} \cup \ldots \cup I_{i+l}^{-}, \quad \text { with } i+1, \ldots, i+l \notin L
$$

By the assumption (8), intervals of this type can exist only if $l \leq k$, and this allows to determine a precise value $\mu^{*}$ such that the construction explained in this section is possible. It is worth noticing that, if $\mu>\mu^{*}$ is fixed, solutions of the boundary value problem (57) give rise to (a restriction of) a function in $\mathcal{N}_{\mu}$ only if $x, y$ are small enough. Indeed, the further condition (C4) has to be satisfied (see Lemma 3.3).

In the next Section 6, we will actually deal (for simplicity of notation) only with problem (38), that is, we will assume that $k=1$. As remarked above, however, the general case $k \geq 1$ could be treated as well.

## 6 Construction of a constrained Palais-Smale sequence

The aim of this section is to construct a (bounded) constrained Palais-Smale sequence $\left(u_{n}\right) \subset \mathcal{N}_{\mu}$ at level (16). This will be done using Ekeland's variational principle [12, Chapter 4] in a careful way. Again (compare with the discussion at the beginning of Section 4) this construction will possibly require to enlarge $\mu^{*}$ more and more times, according to the estimates collected in the previous sections (and ultimately depending on the local estimate contained in Lemma 3.3). In this way, we can still produce a threshold $\mu^{*}$ depending on the weight function $a$ and on the integer $k$, but not on the sequence $\mathcal{L}$ and the integer $N$.

We now start with our arguments. The Ekeland's variational principle applied to $\overline{\mathcal{N}}_{\mu}$ yields the existence of a minimizing sequence $\left(u_{n}\right) \subset \overline{\mathcal{N}}_{\mu}$ for problem (16) such that

$$
\begin{equation*}
J_{\mu}\left(u_{n}\right) \leq J_{\mu}(u)+\frac{1}{n}\left\|u-u_{n}\right\|_{H^{1}\left(I_{N}\right)}, \quad \text { for every } u \in \overline{\mathcal{N}}_{\mu}, n \in \mathbb{N} \tag{58}
\end{equation*}
$$

Our goal now is to show that, for $n$ large enough, $u_{n} \notin \overline{\mathcal{N}_{\mu}} \backslash \mathcal{N}_{\mu}$. Indeed, if this is the case, then (58) implies in a standard way that the constrained gradient $\nabla_{\mathcal{N}_{\mu}} J\left(u_{n}\right) \rightarrow 0$, as required.

We first observe that, in view of Lemma 4.2, the only possibilities for $u_{n}$ to be on $\partial \mathcal{N}_{\mu}$ are given by

$$
\begin{equation*}
\int_{I_{i}^{+}} \dot{u}_{n}^{2}=2\left(c+c_{\zeta}\right), \quad \text { for some } i \in L \tag{59}
\end{equation*}
$$

or, for some $i \in L$,

$$
\left\{\begin{array} { l } 
{ \dot { u } _ { n } ( \sigma _ { i } ^ { - } ) = \rho }  \tag{60}\\
{ u _ { n } ( \sigma _ { i } ) \geq 0 , }
\end{array} \text { or } \left\{\begin{array} { l } 
{ \dot { u } _ { n } ( \sigma _ { i } ^ { - } ) = - \rho } \\
{ u _ { n } ( \sigma _ { i } ) \leq 0 , }
\end{array} \text { or } \left\{\begin{array} { l } 
{ \dot { u } _ { n } ( \tau _ { i } ^ { + } ) = - \rho } \\
{ u _ { n } ( \tau _ { i } ) \geq 0 , }
\end{array} \text { or } \left\{\begin{array}{l}
\dot{u}_{n}\left(\tau_{i}^{+}\right)=\rho \\
u_{n}\left(\tau_{i}\right) \leq 0 .
\end{array}\right.\right.\right.\right.
$$

The rest of this section will be devoted to exclude both the possibilities.
Claim 1. The possibility (59) cannot occur, for n large enough.
Indeed, assume that (59) holds true for some $n$ and $i \in L$; just for simplicity of notation, suppose moreover that $i=0$. Take $p, q$ be as in (31) and let $\bar{u} \in \mathcal{K}_{0}, \bar{u}>0$, according to (4). Then, define $F: \mathbb{R}^{4} \rightarrow \mathbb{R}$ by setting

$$
F(\lambda, \alpha, \beta, \gamma)=\int_{0}^{\tau}\left(\dot{u}^{2}(t ; \lambda, \alpha, \beta)-a^{+} u^{4}(t ; \lambda, \alpha, \beta)\right)-\gamma
$$

where

$$
u(t ; \lambda, \alpha, \beta)=\alpha q(t)+\beta p(t)+\lambda \bar{u}(t)
$$

Notice that $u(0 ; \lambda, \alpha, \beta)=\alpha, u(\tau ; \lambda, \alpha, \beta)=\beta$ and $u(t ; 1,0,0)=\bar{u}(t)$. Since $F(1,0,0,0)=0$ and

$$
\frac{\partial F}{\partial \lambda}(1,0,0,0)=-2 \int_{I_{i}^{+}} \dot{\bar{u}}^{2} \neq 0
$$

the implicit function theorem yields the existence of a function $\lambda=\lambda(\alpha, \beta, \gamma)$, defined a neighborhood of $(0,0,0)$ and having values in a neighborhood of 1 , such that $F(\lambda(\alpha, \beta, \gamma), \alpha, \beta, \gamma)=$ 0 .

In view of Lemma 4.1, we can assume, provided $\mu$ is large enough, that $\alpha_{n}=u_{n}(0)$, $\beta_{n}=u_{n}(\tau)$ and

$$
\gamma_{n}=u_{n}(\tau) \dot{u}_{n}\left(\tau^{+}\right)-u_{n}(0) \dot{u}_{n}\left(0^{-}\right)
$$

are as small as we wish. Then, we can define the function $\widetilde{u}_{n}$ as $\widetilde{u}_{n}(t)=u_{n}(t)$ for $t \notin[0, \tau]$ and

$$
\widetilde{u}_{n}(t)=u\left(t ; \lambda\left(\alpha_{n}, \beta_{n}, \gamma_{n}\right), \alpha_{n}, \beta_{n}, \gamma_{n}\right), \quad \text { for } t \in[0, \tau] .
$$

It is not difficult to see that $\widetilde{u}_{n} \in \overline{\mathcal{N}}_{\mu}$. Moreover, since $\int_{0}^{\tau} \dot{\bar{u}}^{2}=4 c$ and $\widetilde{u}_{n}$ is arbitrarily $H^{1}$-near to $\bar{u}$ for $\mu$ large, we have that

$$
\epsilon_{n}=\frac{1}{4} \int_{0}^{\tau} \dot{\widetilde{u}}_{n}^{2}-c
$$

can be made arbitrarily small. Hence

$$
\begin{aligned}
J_{\mu}\left(u_{n}\right)-J_{\mu}\left(\widetilde{u}_{n}\right) & =\frac{1}{4} \int_{I_{i}^{+}}\left(\dot{u}_{n}^{2}-\dot{\tilde{u}}_{n}^{2}\right)=\frac{c+c_{\zeta}}{2}-\left(c+\epsilon_{n}\right) \\
& \geq \frac{c_{\zeta}-c}{2}-\epsilon_{n} \geq \frac{c_{\zeta}-c}{4} .
\end{aligned}
$$

On the other hand, (58) gives

$$
J_{\mu}\left(u_{n}\right)-J_{\mu}\left(\widetilde{u}_{n}\right) \leq \frac{1}{n}\left\|u_{n}-\widetilde{u}_{n}\right\|_{H^{1}\left(I_{N}\right)}
$$

and this is a contradiction for $n$ large enough (of course, $\left\|u_{n}-\widetilde{u}_{n}\right\|_{H^{1}\left(I_{N}\right)}$ is bounded).

Remark 6.1. The above argument holds also when $c_{\zeta}$ is replaced by any $c^{\prime}>c$. Therefore, using Lemma 4.1 and taking $\mu$ large enough (depending on $c^{\prime}$ ), one can construct minimizing sequences with the property that

$$
\int_{I_{i}^{+}} \dot{u}_{n}^{2}<2\left(c+c^{\prime}\right), \quad \text { for every } i \in L
$$

Claim 2. The possibility (60) cannot occur, for $n$ large enough.
Indeed, assume that (60) holds true for some $n$ and $i \in L$. We will give the details of the proof in the case $u_{n}\left(\tau_{i}\right) \geq 0$ and $\dot{u}_{n}\left(\tau_{i}^{+}\right)=-\rho$ (the other can be treated with similar arguments); moreover, for simplicity of notation we suppose that $i=-1$, that is, we are dealing with the case

$$
u_{n}(\tau-T) \geq 0 \quad \text { and } \quad \dot{u}_{n}\left((\tau-T)^{+}\right)=-\rho .
$$

According to Remark 5.2, we finally assume that $k=1$ and, moreover, we deal with the most difficult case in which the first interval of positivity on the right of $\tau-T$ is not in $L$, that is: $0 \notin L$ and $\pm 1 \in L$. Summing up, we are led to construct a local variation of $u_{n}$ in the interval

$$
I_{-1}^{+} \cup I_{-1}^{-} \cup I_{0}^{+} \cup I_{0}^{-} \cup I_{1}^{+}=[-T, \tau+T] .
$$

Roughly speaking, we will use the variation of Section 5.1 on $I_{-1}^{+}$and $I_{1}^{+}$, and the one of Section 5.2 on the interval $I_{-1}^{-} \cup I_{0}^{+} \cup I_{0}^{-}$. We have to distinguish three situations. In the following, to simplify the notation, we set

$$
x_{n}=u_{n}(\tau-T) \quad \text { and } \quad y_{n}=u_{n}(T) .
$$

If $\left|\dot{u}_{n}\left(T^{-}\right)\right|<\rho$, we argue as follows. We define a family of functions $\widetilde{u}_{n}(\cdot, \xi)$, with $\xi \geq 0$, by setting $\widetilde{u}_{n}(t, \xi)=u_{n}(t)$ for $t \notin[-T, \tau+T]$ and

- for $t \in[-T, \tau-T]$,

$$
\begin{equation*}
\widetilde{u}_{n}(t, \xi)=U\left(u_{n}(-T), x_{n}-\xi, \dot{u}_{n}\left(-T^{-}\right), \dot{\bar{u}}_{\mu}\left((\tau-T)^{+} ; x_{n}-\xi, y_{n}\right)\right)(t) \tag{61}
\end{equation*}
$$

- for $t \in[\tau-T, T]$,

$$
\begin{equation*}
\widetilde{u}_{n}(t, \xi)=\bar{u}_{\mu}\left(t ; x_{n}-\xi, y_{n}\right), \tag{62}
\end{equation*}
$$

- for $t \in[T, \tau+T]$,

$$
\begin{equation*}
\widetilde{u}_{n}(t, \xi)=U\left(y_{n}, u_{n}(\tau+T), \dot{\bar{u}}_{\mu}\left(T^{-} ; x_{n}-\xi, y_{n}\right), \dot{u}_{n}\left((\tau+T)^{+}\right)\right)(t) \tag{63}
\end{equation*}
$$

In formulas (61) and (63), we have denoted (with some abuse of notation) by $U$ a local variation, constructed as in Lemma 5.1, of $u_{n}$ on the interval $[-T, \tau-T]$ and $[T, \tau+T]$, respectively; in (62), $\bar{u}_{\mu}$ is the function constructed in Proposition 5.2. Notice that the map $\xi \mapsto \widetilde{u}_{n}(\cdot ; \xi) \in H^{1}\left(I_{N}\right)$ is of class $C^{1}$ and, of course, $\widetilde{u}_{n}(\cdot, 0)=u_{n}$.

We claim that, for $\xi>0$ small enough, the function $\widetilde{u}_{n}(\cdot, \xi)$ lies in $\overline{\mathcal{N}}_{\mu}$, possibly touching the constraint (60) only outside the intervals we are considering. To show this, the most delicate condition to be checked is the one concerning $\dot{\widetilde{u}}_{n}\left((\tau-T)^{+}, \xi\right)$. We thus observe that, in view of Proposition 5.3,

$$
\left.\frac{d}{d \xi} \dot{\widetilde{u}}_{n}\left((\tau-T)^{+}, \xi\right)\right|_{\xi=0}=-v_{\mu}\left((\tau-T)^{+} ; x_{n}, y_{n}\right)>0
$$

so that $\dot{\tilde{u}}_{n}\left((\tau-T)^{+}, \xi\right)>\dot{u}_{n}\left((\tau-T)^{+}\right)=-\rho$ for small $\xi>0$ as desired.

We now claim that

$$
\begin{equation*}
J_{\mu}\left(\widetilde{u}_{n}(\cdot, \xi)\right) \leq J_{\mu}\left(u_{n}\right)-\left(\frac{\rho}{2}\right) \xi+o(\xi) \quad \text { for } \xi \rightarrow 0^{+} \tag{64}
\end{equation*}
$$

To prove this, we are going to show that

$$
\begin{equation*}
\left.\frac{d}{d \xi} J_{\mu}\left(\widetilde{u}_{n}(\cdot, \xi)\right)\right|_{\xi=0} \leq-\frac{\rho}{2} \tag{65}
\end{equation*}
$$

in the following, we use the notation $J_{\mu,\left[t_{1}, t_{2}\right]}$ for the restriction of the action functional to a subinterval $\left[t_{1}, t_{2}\right] \subset I_{N}$. From (51) of Proposition 5.3, we have

$$
\begin{equation*}
\left.\frac{d}{d \xi} J_{\mu,[\tau-T, T]}\left(\widetilde{u}_{n}(\cdot, \xi)\right)\right|_{\xi=0}=\dot{\bar{u}}_{\mu}\left((\tau-T)^{+} ; x_{n}, y_{n}\right)=\dot{u}_{n}\left((\tau-T)^{+}\right)=-\rho \tag{66}
\end{equation*}
$$

On the other hand, setting $\beta(\xi)=x_{n}+\xi$ and $\beta^{\prime}(\xi)=\dot{\bar{u}}_{\mu}\left((\tau-T)^{+} ; x_{n}-\xi, y_{n}\right)$, we can compute

$$
\begin{aligned}
&\left.\frac{d}{d \xi} J_{\mu,[-T, \tau-T]}\left(\widetilde{u}_{n}(\cdot, \xi)\right)\right|_{\xi=0}=\frac{\partial}{\partial \beta} J_{\mu,[-T, \tau-T]}(U) \frac{d \beta}{d \xi} \quad+\frac{\partial}{\partial \beta^{\prime}} J_{\mu,[-T, \tau-T]}(U) \frac{d \beta^{\prime}}{d \xi} \\
&=\frac{\partial}{\partial \beta} J_{\mu,[-T, \tau-T]}(U) \quad-\frac{\partial}{\partial \beta^{\prime}} J_{\mu,[-T, \tau-T]}(U) \dot{v}_{\mu}\left((\tau-T)^{+} ; x_{n}, y_{n}\right)
\end{aligned}
$$

where we agree that the derivatives with respect to $\beta$ and $\beta^{\prime}$ are evaluated at the point $\left(u_{n}(-T), \dot{u}_{n}(-T), \beta(0), \beta^{\prime}(0)\right)$ while the derivatives with respect to $\xi$ are evaluated at $\xi=0$. From (35) and (36) of Lemma 5.1 we thus obtain

$$
\left|\frac{d}{d \xi} J_{\mu,[-T, \tau-T]}\left(\widetilde{u}_{n}(\cdot, \xi)\right)\right|_{\xi=0} \left\lvert\, \leq \frac{\rho}{8}-C_{\varepsilon} x_{n} \dot{v}_{\mu}\left((\tau-T)^{+} ; x_{n}, y_{n}\right)\right.
$$

From (55) of Lemma 5.3

$$
\left|x_{n} v_{\mu}\left((\tau-T)^{+} ; x_{n}, y_{n}\right)\right| \leq \frac{2 K}{2 T-\tau}+5 \rho
$$

so that, provided $\varepsilon$ is small enough (which, in turns, requires maybe enlarging $\mu^{*}$ )

$$
\left.\left|\frac{d}{d \xi} J_{\mu,[-T, \tau-T]}\left(\widetilde{u}_{n}(\cdot, \xi)\right)\right|_{\xi=0} \right\rvert\, \leq \frac{\rho}{4} .
$$

The arguments for $t \in[T, \tau+T]$ are even simpler, since the dependence of $U$ on $\xi$ is just via $\dot{\bar{u}}_{\mu}\left(T^{-} ; x_{n}-\xi, y_{n}\right)$. Arguing as above, we are thus led to consider the term $C_{\varepsilon} y_{n} v_{\mu}\left(T^{-} ; x_{n}, y_{n}\right)$ which can be estimated using (52) of Lemma 5.3. Hence

$$
\left.\left|\frac{d}{d \xi} J_{\mu,[T, \tau+T]}\left(\widetilde{u}_{n}(\cdot, \xi)\right)\right|_{\xi=0} \right\rvert\, \leq \frac{\rho}{4}
$$

so that, recalling (66), (65) follows.
Finally, we verify that there exists a constant $S$, independent on $n$ (but depending on $\mu$, which however is fixed in this argument), such that

$$
\begin{equation*}
\left\|\widetilde{u}_{n}(\cdot, \xi)-u_{n}\right\|_{H^{1}\left(I_{N}\right)} \leq S \xi+o(\xi) \quad \text { for } \xi \rightarrow 0^{+} \tag{67}
\end{equation*}
$$

To show this, we just need to observe that $\left\|\left.\frac{d}{d \xi} \widetilde{u}_{n}(\cdot, \xi)\right|_{\xi=0}\right\|_{H^{1}\left(I_{N}\right)}$ can be bounded independently on $n$ (we stress once more that the dependence on $\mu$ cannot be avoided, but it is irrelevant for the arguments below, since throughout this section $\mu$ is fixed). This is clear for $t \in[\tau-T, T]$ (since $x_{n}, y_{n} \in[-K, K]$ ) and essentially comes from (34) for $t \in[-T, \tau-T] \cup[T, \tau+T]$ (of course, to bound the derivative $d \widetilde{u}_{n} / d \xi$ one needs - as before - to estimate both $d U / d \beta$ and $d \beta / d \xi$; terms of the type $d \beta / d \xi$ are bounded since $\left.x_{n}, y_{n} \in[-K, K]\right)$.

We are now in position to conclude. Indeed, combining (64) and (67) yields

$$
J_{\mu}\left(\widetilde{u}_{n}(\cdot, \xi)\right)+\frac{1}{n}\left\|\widetilde{u}_{n}(\cdot, \xi)-u_{n}\right\|_{H^{1}\left(I_{N}\right)} \leq J_{\mu}\left(u_{n}\right)+\left(\frac{S}{n}-\frac{\rho}{2}\right) \xi+o(\xi)
$$

and this contradicts (58) for $n$ large and fixed, and $\xi$ sufficiently small. The proof is thus concluded in the case $\left|\dot{u}_{n}\left(T^{-}\right)\right|<\rho$.

If $\dot{u}_{n}\left(T^{-}\right)=-\rho$ (notice that in this case it has to be $u_{n}(T)<0$ ) we argue exactly as before. The only difference is that showing that $\widetilde{u}_{n}(\cdot, \xi) \in \overline{\mathcal{N}}_{\mu}$ now requires the further observation that

$$
\left.\frac{d}{d \xi} \dot{\widetilde{u}}_{n}\left(T^{-}, \xi\right)\right|_{\xi=0}=-v_{\mu}\left(T^{-} ; x_{n}, y_{n}\right)>0
$$

so that, for small $\xi>0, \dot{\widetilde{u}}_{n}\left(T^{-}, \xi\right)>-\rho$ as well.
If $u_{n}\left(T^{-}\right)=\rho$ (notice that in this case it has to be $u_{n}(T)>0$ ) a slightly different argument is needed, since the same variation as before would lead to a function $\widetilde{u}_{n}(\cdot, \xi) \notin \overline{\mathcal{N}}_{\mu}$. Hence, we define here $\widetilde{u}_{n}(t, \xi)=u_{n}(t)$ for $t \notin[-T, \tau+T]$ and

- for $t \in[-T, \tau-T]$,

$$
\widetilde{u}_{n}(t, \xi)=U\left(u_{n}(-T), x_{n}-\xi, \dot{u}_{n}\left(-T^{-}\right), \dot{\bar{u}}_{\mu}\left((\tau-T)^{+} ; x_{n}-\xi, y_{n}-\xi\right)\right)(t)
$$

- for $t \in[\tau-T, T]$,

$$
\widetilde{u}_{n}(t, \xi)=\bar{u}_{\mu}\left(t ; x_{n}-\xi, y_{n}-\xi\right)
$$

- for $t \in[T, \tau+T]$,

$$
\widetilde{u}_{n}(t, \xi)=U\left(y_{n}-\xi, u_{n}(\tau+T), \dot{\bar{u}}_{\mu}\left(T^{-} ; x_{n}-\xi, y_{n}-\xi\right), \dot{u}_{n}\left((\tau+T)^{+}\right)\right)(t)
$$

With this definition, we have that, for $\xi>0$ small enough, the function $\widetilde{u}_{n}(\cdot, \xi)$ lies in $\overline{\mathcal{N}}_{\mu}$. Indeed, from (55) of Lemma 5.3

$$
\left.\frac{d}{d \xi} \dot{\widetilde{u}}_{n}\left((\tau-T)^{+}, \xi\right)\right|_{\xi=0}=-\left(v_{\mu}\left((\tau-T)^{+} ; x_{n}, y_{n}\right)+z_{\mu}\left((\tau-T)^{+} ; x_{n}, y_{n}\right)\right)>0
$$

so that $\dot{\widetilde{u}}_{n}\left((\tau-T)^{+}, \xi\right)>\dot{u}_{n}\left((\tau-T)^{+}\right)=-\rho$ (for small $\left.\xi>0\right)$ and a completely symmetric argument works for $\dot{\widetilde{u}}_{n}\left(T^{-}, \xi\right)$.

At this point, the arguments leading to (64) and (67) are essentially the same as before up to minor modifications and will be omitted.

## 7 Conclusion of the proofs

As for Theorem 2.2, we need to prove the positivity of the solutions found, the $C^{1}$ bound (10), and the properties (P1), (P2), (P3), concerning the behavior for $\mu \rightarrow+\infty$.

Estimate (10) can be easily deduced from Lemma 4.2 and property (C4) in the definition of $\mathcal{N}_{\mu}$ by observing that, because of the convexity/concavity properties of a solution $u,|\dot{u}(t)|$ has local maxima at $t$ if and only $t \in \partial I_{i}^{+}$, with $i \in L$. Lemma 4.2 (points (i) and (ii)) also implies properties (P1) and (P2) (that is, the asymptotic of the solutions on $I_{i}^{-}$and $I_{i}^{+}$with $i \notin L$ ). Property (P3) (that is, the behavior on $I_{i}^{+}$with $i \in L$ ) follows from the following claim: for every $\varepsilon>0$ there exists $\delta>0$ such that, for any $u$ solving $\ddot{u}+a^{+}(t) u^{3}=0$ on $[0, \tau]$ with $r^{2} \leq \int_{0}^{\tau} \dot{u}^{2} \leq 4 c+\delta$ and $u(t) \geq 0$ for $t \in[\zeta, \tau-\zeta]$, the following implication holds:

$$
|u(0)|,|u(\tau)| \leq \delta \quad \Longrightarrow \quad \operatorname{dist}_{W^{2, \infty}}\left(u, \mathcal{K}_{0}\right) \leq \varepsilon
$$

Indeed, by Remark 6.1, we can assume that our solutions satisfy such conditions.
To prove the claim, assume by contradiction that there is $\varepsilon^{*}>0$ and a sequence $\left(u_{n}\right)$ of solutions with $r^{2} \leq \int_{0}^{\tau} \dot{u}_{n}^{2} \rightarrow 4 c, u_{n}(0), u_{n}(\tau) \rightarrow 0$ and

$$
\begin{equation*}
\operatorname{dist}_{W^{2, \infty}}\left(u_{n}, \mathcal{K}_{0}\right)>\varepsilon^{*} . \tag{68}
\end{equation*}
$$

By standard arguments, $u_{n}$ converges in $W^{2, \infty}$ to a limit $\bar{u} \in \mathcal{K}_{0}$, contradicting (68).
We now deal with the positivity, splitting our arguments into three steps.
Step 1. It holds $u\left(\sigma_{i}\right)>0$ and $u\left(\tau_{i}\right)>0$ for every $i \in L$.
To see this, we first recall that $\mathcal{K}_{0}$ is compact. Hence, there exists a constant $m>0, m<\rho$, such that

$$
m \leq \dot{u}\left(\sigma^{+}\right) \quad \text { and } \quad \dot{u}\left(\tau^{-}\right) \leq-m \quad \text { for every } u \in \mathcal{K}_{0}, u>0
$$

Hence, since (P3) holds true, for the solution found in Theorem 2.2 we have that $m / 2 \leq$ $\dot{u}\left(\sigma_{i}^{+}\right)$and $u\left(\tau_{i}^{-}\right) \leq-m / 2$ for every $i \in L$. Convexity arguments now imply the conclusion $\left(u\left(\sigma_{i}\right) \leq 0\right.$ or $u\left(\tau_{i}\right) \leq 0$ would imply a contradiction with (P2)).
Step 2. It holds $u(t)>0$ for every $t \in I_{i}^{+}$with $i \in L$.
This follows from Lemma 3.2.
Step 3. It holds $u(t)>0$ for every $t \in I_{N}$.
This follows from (iii) of Proposition 5.2, since (P1) and (P2) imply that (40) holds true.
We now go back to Theorem 2.1. If $\mathcal{L} \in\{0,1\}^{\mathbb{Z}}$ is a given sequence, the solutions $u=u_{\mathcal{L}, N}$ constructed in Theorem 2.2 for $N$ larger and larger of course are $L^{\infty}$-bounded, independently on $N$. Then, the elementary inequality

$$
\|\dot{v}\|_{L^{\infty}(\mathbb{R})} \leq 2\|v\|_{L^{\infty}(\mathbb{R})}+\|\ddot{v}\|_{L^{\infty}(\mathbb{R})}, \quad \text { for every } v \in W^{2, \infty}(\mathbb{R})
$$

(compare with (24)) implies that $\left\|u_{\mathcal{L}, N}\right\|_{W^{2, \infty}(\mathbb{R})}$ is bounded independently on $N$. Hence, Ascoli-Arzelà Theorem shows that, for $N \rightarrow+\infty, u_{\mathcal{L}, N}$ converges, locally in $W^{2, \infty}$, to a limit function $u_{\mathcal{L}} \in W^{2, \infty}(\mathbb{R})$. It is easily seen that such a function satisfies all the requirements of Theorem 2.1.

Finally, once the proof of Theorem 2.1 is concluded, we have that Theorem 1.1 follows immediately. As far as Theorem 1.2 is concerned, the $W^{2, \infty}(I)$ convergence $u_{\mu}$ to $\bar{u} \in \mathcal{K}_{\mathcal{L}}$ follows from properties (P2) and (P3) if $I=I_{i}^{+}$, and from Lemma 3.4 if $I=\left[\tau_{i}+\delta, \sigma_{i+1}-\delta\right]$, for any $\left[\tau_{i}, \sigma_{i+1}\right]=I_{i}^{-}$and $\delta>0$. The $H_{\mathrm{loc}}^{1}$ convergence descends from the $W^{2, \infty}$ one,
together with (P1). To prove the Hölder convergence, we first observe that (P1), (P2), (P3) also imply that $u_{\mu} \rightarrow \bar{u}$ uniformly. On the other hand, since for every $s \neq t$

$$
\text { both } \quad \frac{\left|u_{\mu}(t)-u_{\mu}(s)\right|}{|t-s|} \leq \rho \quad \text { and } \quad \frac{|\bar{u}(t)-\bar{u}(s)|}{|t-s|} \leq \rho
$$

(recall Remark 5.1), we have

$$
\frac{\left|\left(u_{\mu}-\bar{u}\right)(t)-\left(u_{\mu}-\bar{u}\right)(s)\right|}{|t-s|^{\alpha}} \leq(2 \rho)^{\alpha}\left|\left(u_{\mu}-\bar{u}\right)(t)-\left(u_{\mu}-\bar{u}\right)(s)\right|^{1-\alpha} \rightarrow 0
$$

by uniform convergence, concluding the proof.
Remark 7.1. As we mentioned, our method can be applied also in different situations. To conclude, we briefly discuss some of them, with an emphasis on the minor changes they require to be dealt with.

Changing sign solutions. As $\mathcal{K}_{0}$ is defined by even conditions, it contains both positive and negative solutions of $\ddot{u}+a^{+}(t) u^{3}=0$ in $H_{0}^{1}(0, \tau)$. As a consequence, if $\mathcal{L} \in\{-1,0,1\}^{\mathbb{Z}}$, one can redefine $\mathcal{K}_{\mathcal{L}}$ as

$$
\mathcal{K}_{\mathcal{L}}=\left\{u \in W^{1, \infty}(\mathbb{R}):\left.u\right|_{I_{i}^{+}} \in \mathcal{K}_{i}, \pm\left. u\right|_{I_{i}^{+}}>0 \text { if } \mathcal{L}_{i}= \pm 1 \text { and } u \equiv 0 \text { elsewhere }\right\}
$$

and obtain changing sign solutions by singularly perturbing such a set. In doing that, one has to change the definition of $\mathcal{N}_{\mu}$ (equation (15)), by substituting property (C2) with
(C2') $\pm u(t)>0$ for every $t \in\left[\sigma_{i}+\zeta, \tau_{i}-\zeta\right]$ with $\mathcal{L}_{i}= \pm 1$.
Consequently, one can find solutions having the appropriate sign on each $I_{i}^{+}$, when $\mathcal{L}_{i}= \pm 1$, and changing sign exactly once in every connected components of the complementary set.

Periodic weights with more than two nodal intervals. If $a$ is periodic, but changes sign more than once in $[0, T]$ (but still a finite number of times), then the main problems arise from a notational point of view, while the conceptual one should be clear: to start with, one has to partition $[0, T]$ in the juxtaposition of consecutive intervals $I_{i}^{ \pm}$, according to the rule that, for every $I_{i}^{-}=\left[\tau_{i}, \sigma_{i+1}\right]$ and $\delta>0$ small, it holds

$$
\int_{\tau_{i}}^{\tau_{i}+\delta} a^{-}(t) d t>0 \quad \text { and } \quad \int_{\sigma_{i+1}-\delta}^{\sigma_{i+1}} a^{-}(t) d t>0
$$

(notice that this is always possible, exactly as in the one-zero case, up to carefully choosing the points $\left.\sigma_{i}, \tau_{i}\right)$. Accordingly, $\mathcal{N}_{\mu}$ should be defined by means of different (still finite) constants $r_{i}, \rho_{i}, \zeta_{i}$. As a byproduct, this should prove multiplicity of periodic solutions.

Non periodic weights. In the same spirit of what we have just enlightened, one may also consider non necessarily periodic weights $a$, though enjoying some uniform oscillatory properties (in particular, the constants $r_{i}, \rho_{i},\left|I_{i}^{ \pm}\right|, i \in \mathbb{Z}$, should be bounded below and above independently on $i$; furthermore, a uniform version of Lemma 3.2 should hold). By considering periodic truncation of such a weight on larger and larger intervals, one should eventually obtain bounded entire solutions, with possible recursivity properties.

Dirichlet and Neumann boundary conditions. Beyond periodic ones, other boundary conditions on bounded intervals $I$ can be considered in our construction. Dirichlet homogeneous ones can be obtained by replacing $H_{\mathrm{per}}^{1}(I)$ with $H_{0}^{1}(I)$, thus recovering the results by Gaudenzi, Habets and Zanolin [14, 16]. In the same way, using $H^{1}(I)$, one can solve the Neumann homogeneous problem on $I$. From this point of view, the Neumann problem, as well as the periodic one, contains the further difficulty that the quadratic part of the action functional does not correspond to an equivalent norm in the ambient Hilbert space.

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